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## Original Article

# First-principles calculations to investigate HgY<sub>2</sub>S/Se<sub>4</sub> spinel chalcogenides for optoelectronic and thermoelectric applications



Ghulam M. Mustafa <sup>a,\*</sup>, Sadaf Saba <sup>b</sup>, N.A. Noor <sup>c</sup>, A. Laref <sup>d</sup>,  
Magda Abd El-Rahman <sup>e,f</sup>, Zahid Farooq <sup>a</sup>, R.B. Behram <sup>g</sup>, Zaka Ullah <sup>a</sup>

<sup>a</sup> Department of Physics, Division of Science and Technology, [University of Education](http://www.ue.edu.pk), Lahore, Punjab 54770, Pakistan

<sup>b</sup> Computational Materials Modeling Laboratory, Department of Physics, Government College University, Faisalabad, 38040, Pakistan

<sup>c</sup> Department of Physics, RIPHAH International University, Campus Lahore, Pakistan

<sup>d</sup> Department of Physics and Astronomy, College of Science, King Saud University, Riyadh, 11451, Saudi Arabia

<sup>e</sup> Department of Physics, College of Science, King Khalid University, Abha, 61413, Saudi Arabia

<sup>f</sup> Department of Radiation Physics, National Center of Radiation Research and Technology (NCRRT), Atomic Energy Authority, 11787Cairo, Egypt

<sup>g</sup> Department of Physics, AIU, Islamabad, Pakistan

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## ABSTRACT

Spinel chalcogenides have great potential for optoelectronic and thermoelectric applications and therefore received huge attention in recent years. In this regard, here we investigate the structural, optical, electronic, and thermoelectric characteristics of HgY<sub>2</sub>S<sub>4</sub> and HgY<sub>2</sub>Se<sub>4</sub> spinel chalcogens using a density functional theory-based WIEN2k package. Ground state optimization of crystal structure and thermodynamic stability of the material is probed from the energy volume optimization graph and computation of enthalpy of formation. The computation of Poisson's and Pugh's ratios revealed the ductile nature of these materials. Bandgap calculation is performed using TB-mBJ package, which exposed the direct band nature of these semiconducting materials with bandgap values of 1.2 eV for HgY<sub>2</sub>S<sub>4</sub> and 0.6 eV for HgY<sub>2</sub>Se<sub>4</sub>. The evaluation of optical characteristics and transport features revealed these compositions' potential for optoelectronic and thermoelectric applications.

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\* Corresponding author.

E-mail address: [dr.ghulam.muhammad@ue.edu.pk](mailto:dr.ghulam.muhammad@ue.edu.pk) (G.M. Mustafa).

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## 1. Introduction

Thermoelectric materials are the potential alternatives for harvesting devices like thermoelectric generators whose efficiency is measured in the form of the figure of merit (ZT). This figure of merit depends upon the Seebeck coefficient ( $S$ ), electrical conductivity ( $\sigma$ ), and thermal conductivity ( $\kappa$ ). The value of ZT can be calculated using the expression  $ZT = S^2\sigma T/\kappa$  [1] where the term  $S^2\sigma$  is called the power factor (PF) and determines the efficiency of thermoelectric materials without considering the effect of thermal conductivity [2]. The material with a high value of PF is usually heavily doped semiconductors like  $\text{Bi}_2\text{Te}_3$ . In semiconductors, the thermal conductivity mainly comes from the lattice thermal conductivity governed by phonons [3]. For better performance of any thermoelectric material it must have higher values of  $S$ ,  $\sigma$ , and PF with minimal  $\kappa$  [4]. Thermoelectric materials with a small thermal conductivity value are highly desirable for advanced technological applications [5]. In this paper, we evaluated the spinel structured material with the general formula  $\text{AM}_2\text{X}_4$  where A is divalent, M is a trivalent atom, and X is a chalcogen atom. Atom A occupies 1/8th of the tetrahedral positions, and atom M occupies half of the octahedral positions, whereas X atoms form the FCC lattice [6]. Spinel compositions cover a broad spectrum of applications because of the significant mixing of cations at A and M sites ranging from optical to electronic, magnetic to optoelectronic, and medical to industrial. For their potential applicability in thermoelectric applications, Spitzer proposed that the high coordination number of M atoms favors reducing the thermal conductivity of spinels [7]. High concentrations of octahedral holes present in the lattice act as a scattering center for phonons and thus reduce thermal conductivity. Therefore, this kind of composition with a multi-valley electronic structure is supposed to reveal a high value of power factor [8].

Spinel chalcogenides have been extensively studied compositions in the last few decades. More than 300 known spinel compositions are reported with  $X = \text{S}$  or  $\text{Se}$  in which most have 3d transition metals at M and A-sites [9]. In addition to spinel chalcogenides, oxides-based spinel compositions have also been explored to evaluate their magnetic response, but they are not feasible for thermoelectric devices because of their very high bandgap values [6]. Furthermore, the sulfides-based compositions are usually conductors or small band semiconductors, for example, sulfides of Fe, Ni, Co, and V are not stable in spinel and transform to  $\text{M}_3\text{S}_4$  type under high pressure and temperature. The spinels with  $X = \text{Te}$  are conductors with a small value of power factor. Spinel selenides are exceptional thermoelectric materials because of their low thermal conductivity and their wide spectrum of electronic properties ranging from metals to insulators [10]. Therefore, optimizing doping concentration to obtain minimal thermal power is desirable but challenging for  $p$  and  $n$ -type semiconductors [11]. The Skutterudites and  $\text{Zn}_4\text{Sb}_3$  are proven thermoelectric materials but could not accommodate a variety of dopants to obtain minimum thermal conductivity. Thanks to the chemical versatility of  $\text{A}_x\text{M}_{3-x}\text{X}_4$  compositions, these compositions can accommodate a variety of dopants which allows tuning the thermal conductivity and other

thermoelectric properties [12]. Spinel chalcogenides are promising materials and get huge attention because of their broad spectrum applications in thermoelectric devices [13]. Here, we evaluated the structural, optical, electronic, and thermoelectric characteristics of  $\text{HgY}_2\text{S}_4$  and  $\text{HgY}_2\text{Se}_4$  compounds by first-principles calculations. Because, Y-based spinel chalcogenides have been studied as cost-effective magnetic materials, so they have great potential for spintronic applications like quantum computing, quantum valve, and magnetic memory devices [14]. The presence of Y in the spinel lattice provides a conducting channel for Hg ions and make them good for low-temperature catalytic property [15]. In addition to spinels, a detailed investigation of electronic, optical, mechanical, and thermodynamic studies of full, half, and quaternary Heusler alloys [16–21]. On the other hand, double perovskite materials have also been extensively explored for their optoelectronic and thermoelectric applications for futuristic devices [22–27].

However, a limited number of investigations are available to understand the mercury adsorption and oxidation method of S and Se-based chalcogenides. Since there are no theoretical or experimental reports in the literature available for thermoelectric investigations of Y-based spinels  $\text{HgY}_2\text{S}/\text{Se}_4$ , and here we first time explored the optical, electronic, and thermoelectric properties of these materials to check their potential for optoelectronic and thermoelectric applications. The higher values of the figure of merit (ZT) and the computed small bandgap indicate that Y-based spinels  $\text{HgY}_2\text{S}/\text{Se}_4$  are suitable for energy conversion applications.

## 2. Computational analysis

The optoelectronic and thermoelectric properties of Y-based spinels  $\text{HgY}_2\text{S}/\text{Se}_4$  were computed within the framework of the WIEN2k code [28–33], based on density functional theory. We employed the generalized gradient approximation (GGA) proposed by Perdew, Burke, and Ernzerhof (PBE) for the electronic exchange-correlation energy functional. The potentials were created by considering the outermost electrons of an atom and give accurate results of structure parameters comparable to experiment. For exchange-correlation energy functional, local density approximation (LDA) is also used, but it give results underestimate as compared to GGA. We did modified Becke-Johnson (mBJ) potential calculations by considering electronic and optical properties and found that mBJ potential slightly increase the bandgap [34–40]. Furthermore, modified Becke-Johnson potential (mBJ) has been employed for calculating the bandgap close to the actual band gap of the material because this code generates more reliable bandgap values [41]. A variety of approximations have been utilized to compute different physical features of these materials [42–50]. The ground state structural parameters are calculated by relaxing and optimizing the unit cell which leads to estimate the electrical and optical characteristics of the  $\text{HgY}_2\text{S}/\text{Se}_4$ . In the interstitial region, the cut-off parameter of the plane wave function was set. The order of k-mesh was  $10 \times 10 \times 10$ , which generated 1000 k-points. On the other hand, the values of  $R_{\text{MT}} \times K_{\text{max}}$  and  $G_{\text{max}}$  were kept at 8 and 16, respectively where  $R_{\text{MT}}$  represent the radius of smallest unit

cell in reciprocal space called the muffin tin radius and  $K_{\max}$  denoted the maximum value of K-vector. The energy was converged as  $10^{-5}$  Ry. Afterward, using the BoltzTrap code [51], the thermoelectric parameters of  $\text{HgY}_2\text{S}/\text{Se}_4$  spinels were investigated.

### 3. Results and discussion

#### 3.1. Structural parameters

Fig. 1 presents the crystallographic arrangement of  $\text{HgY}_2\text{S}/\text{Se}_4$  spinels in stick and ball format and then in a polyhedral form. The blue-colored cations are of Hg placed at the tetrahedral site, black-colored cations are of Y placed at the octahedral site, and yellow-colored anions are of S/Se developing FCC lattices [52].

Before calculating any physical property, the structure was first relaxed in its ground state to determine the ground state lattice parameters. The fitting of structural data with the Birch-Murnaghan equation of state generated the value of ground state optimized volume of unit cell which is used to compute the lattice constant  $a_0(\text{\AA})$  and Bulk modulus  $B_0(\text{GPa})$ . The computed values of  $a_0$  and  $B_0$  are presented in Table 1. It can be noticed from these values that the value of  $a_0$  enhanced while the value of  $B_0$  reduced by keeping S at Se place. This variation in  $a_0$  and  $B_0$  is caused by the higher ionic radius of Se rather than S. The thermodynamic stability is further verified by measuring the enthalpy of formation ( $\Delta H_f$ ).

$$\Delta H_f = E_{\text{total}}(\text{Hg}_l\text{Y}_m\text{S}/\text{Se}_n) - lE_{\text{Hg}} - mE_{\text{Y}} - nE_{\text{S/Se}} \quad (1)$$

$E_{\text{total}}(\text{Hg}_l\text{Y}_m\text{S}/\text{Se}_n)$  is the overall energy, whereas  $E_{\text{Hg}}$ ,  $E_{\text{Y}}$ , and  $E_{\text{S/Se}}$  are the energies of Hg, Y, and S/Se, respectively.  $\text{HgY}_2\text{S}/\text{Se}_4$  spinels are thermodynamically stable because the measured values of  $\Delta H_f$  have negative energy values.

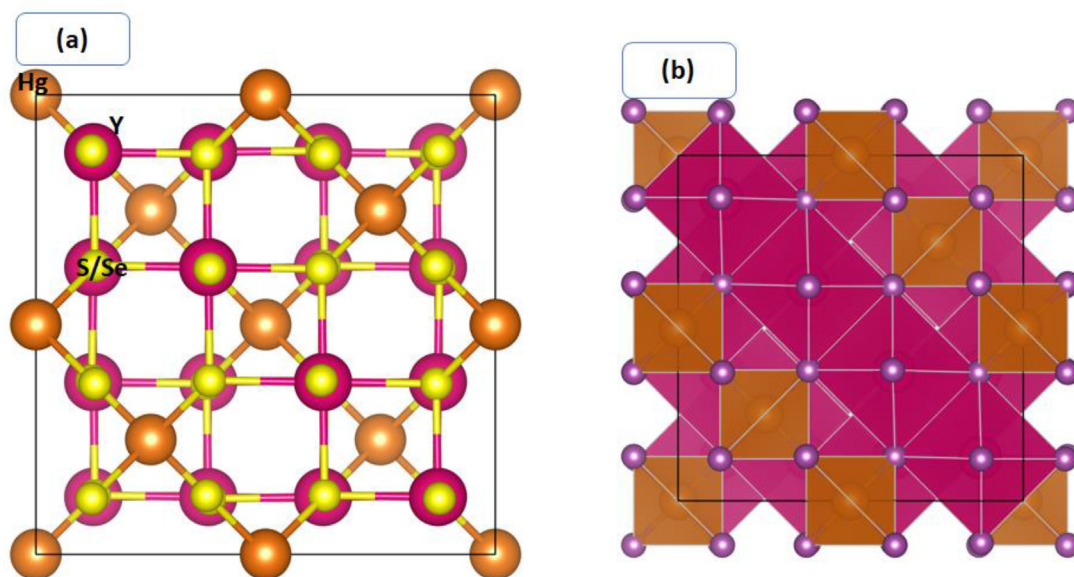
The mechanical response of  $\text{HgY}_2\text{S}/\text{Se}_4$  spinels was estimated by computing three elastic constants i.e.,  $C_{11}$ ,  $C_{12}$ , and

$C_{44}$  which pave the path to calculate various moduli of elasticity. These elastic constants help to define the stability criteria for cubic systems ( $C_{11} > 0$ ,  $C_{11} + 2C_{12} > 0$ , and  $C_{11} - C_{12} > 0$ ) called the Born stability criteria. The computed values of these constants also support the stability of these compositions. Using the calculated values of elastic constant, we calculate bulk modulus (B), Young modulus (Y), and Shear modulus (S), and their estimated values are given in Table 1. The measured values of bulk modulus  $B_0$  (GPa) from elastic parameters and computed from structural optimization data represent that studied spinels reveal better mechanical properties. To identify whether the material is ductile or brittle, the Pugh ratio (B/G) and Poisson ratio ( $\nu$ ) play their key role because if the value of B/G is  $\geq 1.75$  the material is ductile other wise brittle similarly the if the value of  $\nu \geq 0.26$  the material is ductile other brittle. In present scenario, the value of B/G is 1.75 and 1.78 where as value of  $\nu$  is 0.26 and 0.27 for  $\text{HgY}_2\text{S}_4$  and  $\text{HgY}_2\text{Se}_4$  respectively which assured their ductile nature [53].

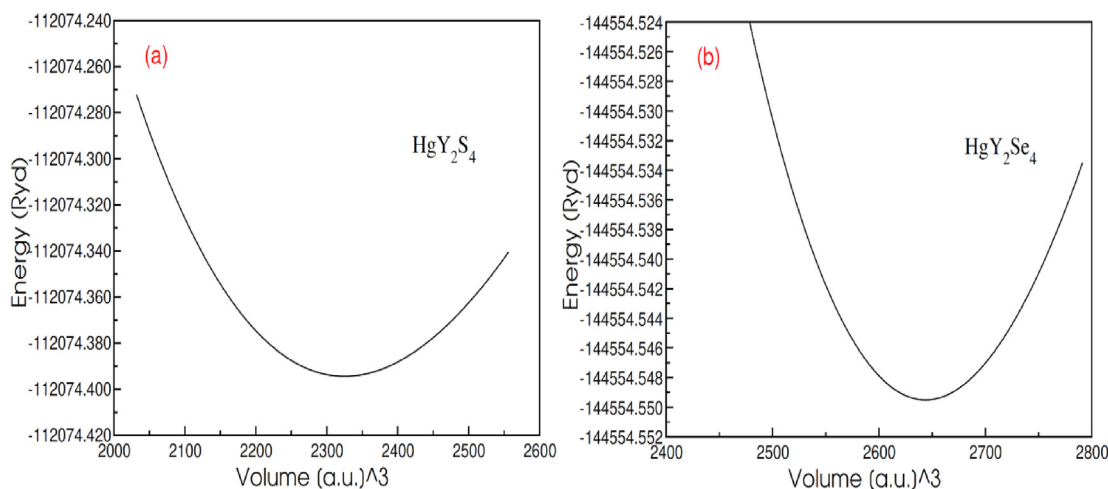
Fig. 2(a) and (b) present the volume energy optimization graphs for  $\text{HgY}_2\text{S}_4$  and  $\text{HgY}_2\text{Se}_4$ . The value of the optimized

**Table 1 – Computed structural and elastic parameters of  $\text{HgY}_2\text{S}_4$  and  $\text{HgY}_2\text{Se}_4$  spinels.**

| Parameters        | $\text{HgY}_2\text{S}_4$ | $\text{HgY}_2\text{Se}_4$ |
|-------------------|--------------------------|---------------------------|
|                   | PBEsol-GGA               | PBEsol-GGA                |
| $a_0(\text{\AA})$ | 11.12                    | 11.62                     |
| $B_0$ (GPa)       | 76.19                    | 63.65                     |
| $\Delta H_f$      | −8.64                    | −6.88                     |
| $C_{11}$          | 150.23                   | 112.40                    |
| $C_{12}$          | 39.84                    | 32.68                     |
| $C_{44}$          | 37.45                    | 29.85                     |
| B                 | 73.63                    | 59.25                     |
| G                 | 42.05                    | 33.49                     |
| Y                 | 110.29                   | 84.56                     |
| B/G               | 1.75                     | 1.77                      |
| $\nu$             | 0.26                     | 0.27                      |
| A                 | 0.68                     | 0.75                      |



**Fig. 1 – Unit cell of  $\text{HgY}_2\text{S}/\text{Se}_4$  in (a) ball-stick model and (b) polyhedral format.**



**Fig. 2 – Energy-volume optimization graph of  $\text{HgY}_2\text{S}_4$  and  $\text{HgY}_2\text{Se}_4$  spinels.**

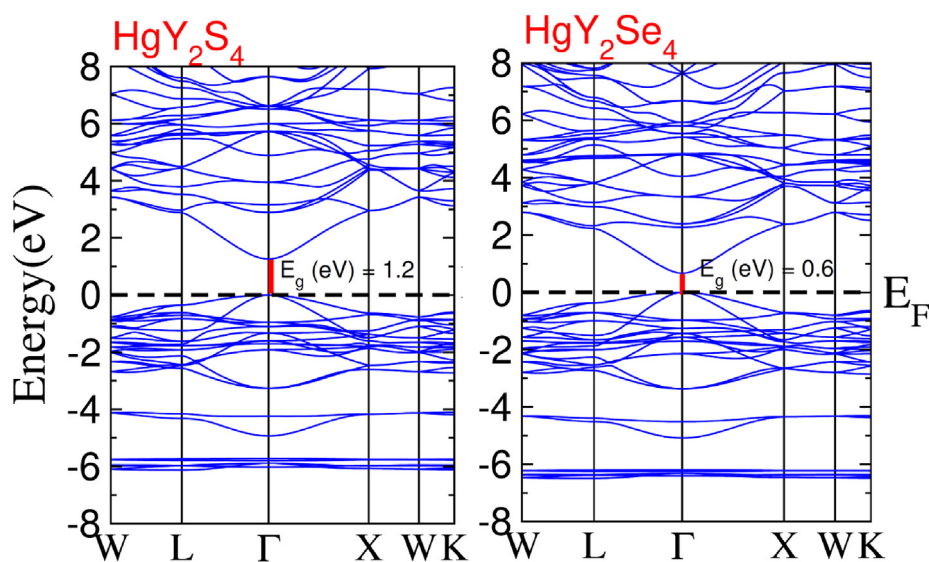
lattice constant is determined from the lowest point on this curve. By keeping S at Se place, the lowest point appeared at a relatively higher volume value, indicating the unit cell's expansion. The larger ionic radii of Se, contrary to the S expanded the unit cell.

### 3.2. Electronic band structure

The band structure analysis aims to understand the material's nature and appropriateness in device fabrications. The band structure of  $\text{HgY}_2\text{S}/\text{Se}_4$  spinels is observed in Fig. 3. Both compositions possess direct band gap structure and small bandgap values for both compositions reveal their semi-conducting nature. The computed bandgap ( $E_g$ ) value for  $\text{HgY}_2\text{S}_4$  is found to be 1.2 eV which is reduced to 0.6 eV when S is replaced with Se. The higher atomic radii of Se caused this reduction in  $E_g$  value compared to S atoms [54–56]. These narrow bandgap semiconductors are appropriate for

optoelectronic devices working in the infrared region of the electromagnetic spectrum.

Interband jumping probability of electrons is governed by density of states and band structure. This is because the hybridization of involving states plays a decisive role in defining the bandgap and, thus, the electronic transition. To have a deep look of involving states of contributing atoms in band formation, the graphs of total density of states (TDOS) and partial density of states (PDOS) are presented in Fig. 4. A closer look at the PDOS revealed that 5d-states of Hg, 4s-states of Y and 3p-states of S and Se form the valence band in which 5d-states of Hg are deep in the valence band, appeared around  $-5.5$  eV for  $\text{HgY}_2\text{S}_4$ , and move further down to  $-6$  eV for  $\text{HgY}_2\text{Se}_4$ . However, 4s-states of Y spread between  $-3.5$  and  $-0.5$  eV. However, the valence band edge is formed by the 3p-states of S and Se, which touches the Fermi level. The conduction band is composed of 4s-states of Y and 3p states of S and Se. When Se replaces S, 4s states of Y and 3p states of Se in



**Fig. 3 – Band structure of  $\text{HgY}_2\text{S}_4$  and  $\text{HgY}_2\text{Se}_4$  compositions.**

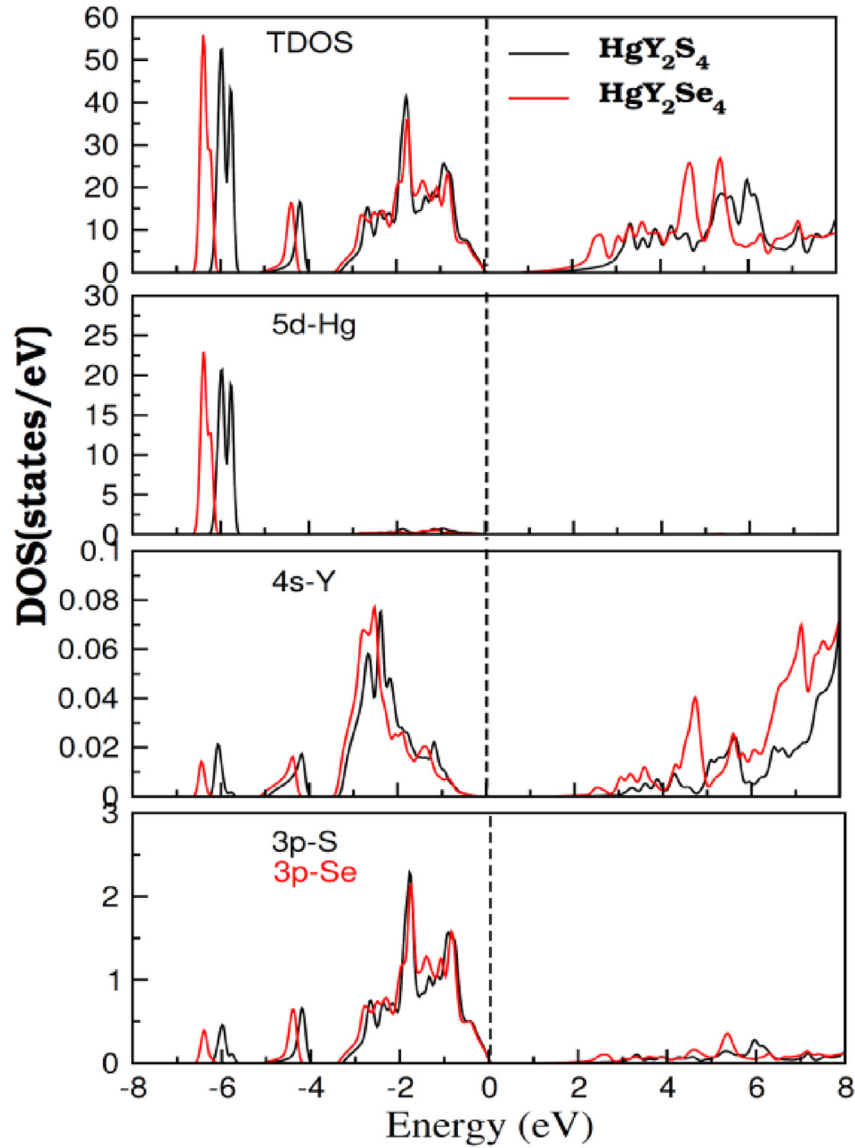


Fig. 4 – Total and partial density of states for  $\text{HgY}_2\text{S}_4$  and  $\text{HgY}_2\text{Se}_4$  spinels.

the CB moved towards the fermi level, which caused the net reduction in the  $E_g$  value from 1.2 to 0.6 eV. This replacement is a powerful approach to tuning the spinel chalcogenides bandgap, making them suitable for optoelectronic and thermoelectric devices [57,58].

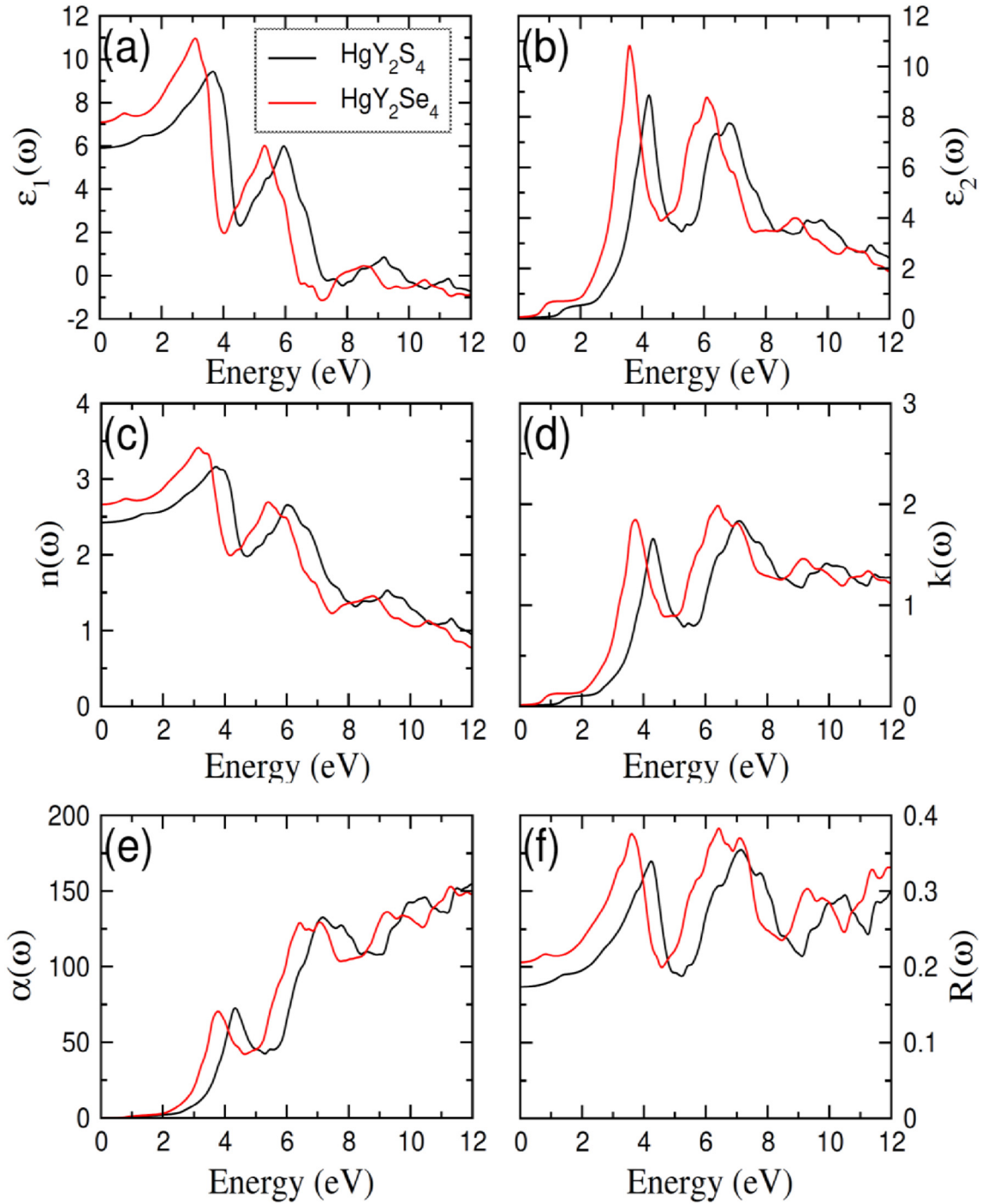
### 3.3. Optical characteristics

The calculated values of band gap energy ( $E_g$ ) and static dielectric constant  $\epsilon_1(0)$  are given in Table 2 for  $\text{HgY}_2\text{S}_4$  and  $\text{HgY}_2\text{Se}_4$  compositions along with refractive index ( $n(0)$ ) and reflectivity ( $R(0)$ ) values. It is observed that values of  $\epsilon_1(0)$ ,  $n(0)$ , and  $R(0)$  are increased by replacing S atoms with Se. The Se atoms have higher atomic radii, so their density is high that's why the refractive index is high. Similarly, its polarizability is higher than  $\text{HgY}_2\text{S}_4$  composition; therefore, it shows a higher value of  $\epsilon_1(0)$  [59].

Table 2 – Energy bandgap  $E_g$ (eV), real part of dielectric constant  $\epsilon_1(0)$ , refractive index  $n(0)$  and reflectivity  $R(0)$  at 0 energy for  $\text{HgY}_2\text{S}_4$  and  $\text{HgY}_2\text{Se}_4$ .

| Parameters      | $\text{HgY}_2\text{S}_4$ | $\text{HgY}_2\text{Se}_4$ |
|-----------------|--------------------------|---------------------------|
|                 | Our work                 |                           |
| $E_g$ (eV)      | 1.20                     | 0.60                      |
| $\epsilon_1(0)$ | 5.89                     | 7.09                      |
| $n(0)$          | 2.42                     | 2.66                      |
| $R(0)$          | 0.17                     | 0.21                      |

The dielectric constant is a complex quantity  $\epsilon(\omega) = \epsilon_1(\omega) + i\epsilon_2(\omega)$ , whose real  $\epsilon_1(\omega)$  and imaginary  $\epsilon_2(\omega)$  parts were estimated to understand better the optical response of  $\text{HgY}_2\text{S}_4$  and  $\text{HgY}_2\text{Se}_4$  spinels [60]. The dispersion and absorption of light can be observed when there is an interaction of photons of light with matter. The former term in  $\epsilon(\omega)$



**Fig. 5 – Energy dependent (a)  $\epsilon_1(\omega)$ , (b)  $\epsilon_2(\omega)$ , (c)  $n(\omega)$ , (d)  $k(\omega)$ , (e)  $\alpha(\omega)$ , and (f)  $R(\omega)$  for  $\text{HgY}_2\text{S}_4$  and  $\text{HgY}_2\text{Se}_4$  spinels.**

corresponds to the absorption of light, while the latter one corresponds to the dispersion of light. The separate frequency dependent response of real and imaginary part of dielectric constant can be understood based on the Kramers-Kronig relation [61]:

In Fig. 5(a–d), the plots of  $\epsilon_1(\omega)$ ,  $\epsilon_2(\omega)$ ,  $n(\omega)$  and  $k(\omega)$  are shown where it was noticed that the highest peak of  $\epsilon_1(\omega)$  was noticed at 3.8 eV for  $\text{HgY}_2\text{S}_4$  and 3 eV for  $\text{HgY}_2\text{Se}_4$ . Since the

energy and frequency of electromagnetic waves are directly linked, therefore moving of highest intensity peak towards lower energy causes an increase in  $\epsilon_1(\omega)$  value from 9 to 11 [62]. In Fig. 5(b), the highest peak of  $\epsilon_2(\omega)$  was noticed at 4.1 and 3.8 eV for  $\text{HgY}_2\text{S}_4$  and  $\text{HgY}_2\text{Se}_4$ , respectively. For the former composition, the value of  $\epsilon_2(\omega)$  is zero below 1 eV and subsequent composition this energy limit is 0.6 eV. These spinels are appropriate candidates for optoelectronic

applications in the infrared region. Moreover,  $\epsilon_1(\omega)$  is related to the  $n$  and  $k$  through  $n^2 + k^2 = \epsilon_1(\omega)$  relation. The term  $n(\omega)$  determines how often light slows down in a medium than in a vacuum. The highest intensity peaks for  $n(\omega)$  and  $k(\omega)$  are seen at 3.9 eV and 7 eV, respectively for  $\text{HgY}_2\text{S}_4$  and these peaks are shifted to 3.1 eV and 6.3 eV, for  $\text{HgY}_2\text{Se}_4$ . To further understand the absorption of light, the absorption coefficient  $\alpha(\omega)$  was calculated and depicted in Fig. 5(e) [63]. In addition, the value of  $\alpha(\omega)$  quantifies the absorbed light by semiconductors and it became significant when incident light is smaller than bandgap (CB) value. Its value is zero below 2 eV and has maximum peaks at 12 eV for both compositions. It means the studied spinels will absorb more light for higher energies. Absorption, transmission, and reflection all happen concurrently as light strike a surface. Another essential factor for optoelectronic applications, reflectivity  $R(\omega)$ , is shown in Fig. 5(f), which is used to examine the light that is reflected from the surface. In contrast with  $\alpha(\omega)$ , the values of  $R(\omega)$  are in fractions that's why the performance of optical devices will not be affected [64].

### 3.4. Thermoelectric characteristics

For device applications, it is much needed to measure the thermoelectric factors of the materials because the ability of the device to alter heat into electrical energy can be calculated through these factors [65]. In Fig. 6(a–d), the thermoelectric parameters are plotted in the temperature range of 200–800 K, and their measured readings are given in Table 3.

The moving free carriers correspond to the  $\sigma/\tau$ , and its value gradually increased with an increase in temperature for  $\text{HgY}_2\text{S}_4$  and  $\text{HgY}_2\text{Se}_4$  spinels [66]. The carriers transition from the valence to the conduction band after getting energy from temperature. The values of  $\sigma/\tau$  are lower at lower temperatures and maximum at 800 K temperature for both compositions. Similarly, the heat caused by the carriers and lattice vibrations corresponds to the  $\kappa/\tau$ . These lattice vibrations generate the phonon waves. The trend of thermal conductivity concerning temperature is almost similar to that of electrical conductivity, which increases with increasing temperature, but its values are much lower than electrical conductivity [67]. For thermoelectrics, the studied compositions are appropriate candidates owing to their small value of thermal conductivity ( $10^{14}$ ) and high value of electrical conductivity ( $10^{19}$ ). The voltage due to the temperature gradient (Seebeck coefficient ( $S$ )) is graphed concerning the temperature in Fig. 6(b). The graph shows that the value of  $S$  is low for  $\text{HgY}_2\text{S}_4$  and high for  $\text{HgY}_2\text{Se}_4$ . The positive value of  $S$  reveals that these compositions have a p-type semiconducting response [68]. From Fig. 6(d), the plot of the power factor (PF) can be seen through which the thermoelectric efficiency of the material can be found. An increasing trend in PF is noticed for  $\text{HgY}_2\text{S}_4$  and  $\text{HgY}_2\text{Se}_4$  spinels against temperature. This factor is calculated to understand the thermoelectric performance of the materials.

The outcome of the combined thermoelectric response is sensed from the figure of merit (ZT). Fig. 7 shows the change of ZT versus temperature, where it increased with growing

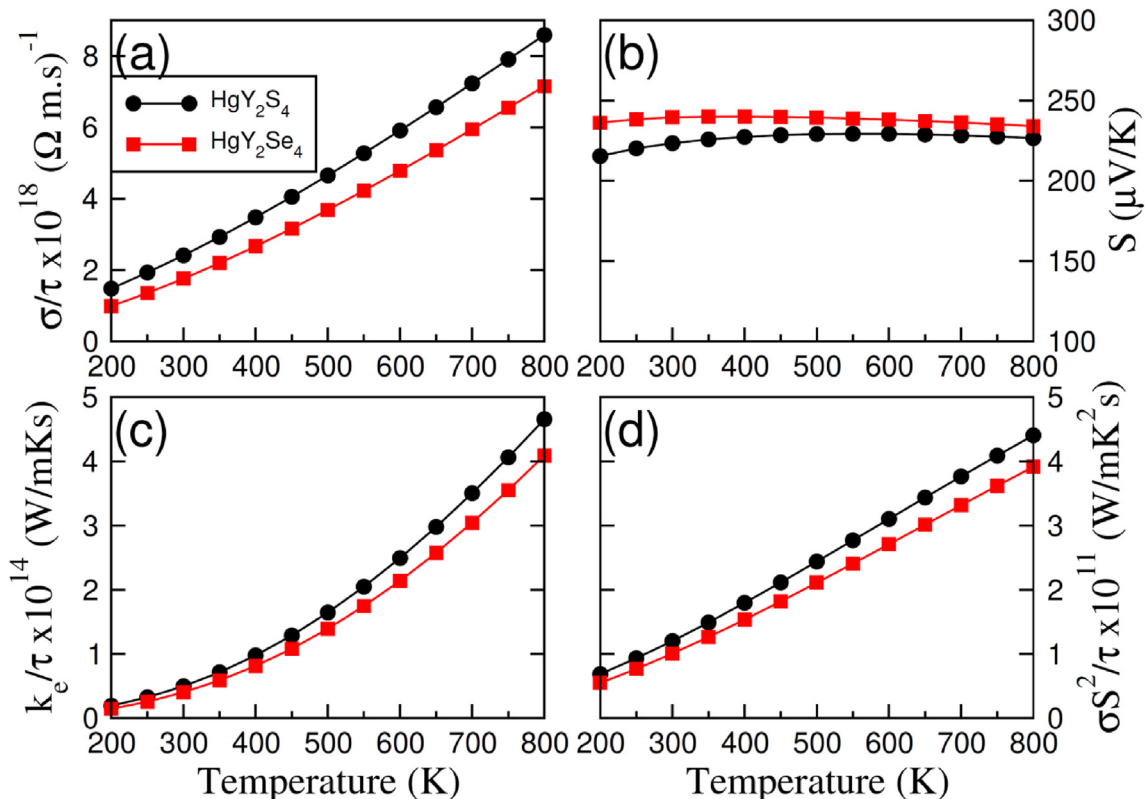
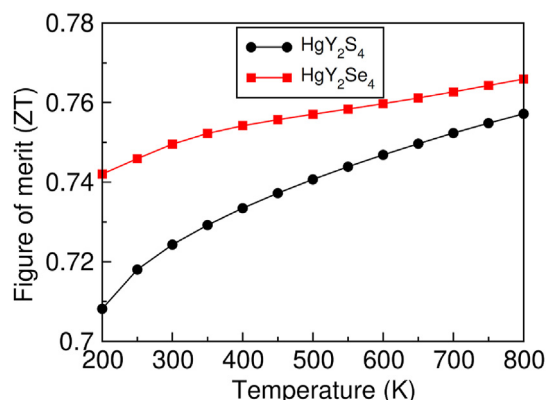


Fig. 6 – Temperature dependent (a) Electrical conductivity ( $\sigma/\tau$ ), (b) Seebeck coefficient ( $S$ ), (c) thermal ( $\kappa_e/\tau$ ) conductivity, and (d) power factor ( $\sigma S^2/\tau$ ) of  $\text{HgY}_2\text{S}_4$  and  $\text{HgY}_2\text{Se}_4$ .

**Table 3 – Thermoelectric parameters of  $\text{HgY}_2\text{S}_4$  and  $\text{HgY}_2\text{Se}_4$  at RT.**

| Spinel                    | $\sigma/\tau$ ( $\times 10^{19}/\Omega\text{ms}$ ) | $\kappa_e/\tau$ ( $\times 10^{14}\text{W/mK}$ ) | $S(\mu\text{V/K})$ | $\sigma S^2/\tau$ | ZT   |
|---------------------------|----------------------------------------------------|-------------------------------------------------|--------------------|-------------------|------|
| $\text{HgY}_2\text{S}_4$  | 2.41                                               | 0.49                                            | 223.36             | 1.20              | 0.72 |
| $\text{HgY}_2\text{Se}_4$ | 1.76                                               | 0.40                                            | 239.48             | 1.01              | 0.75 |

**Fig. 7 – Temperature-dependent figure of merit (ZT) of  $\text{HgY}_2\text{S}_4$  and  $\text{HgY}_2\text{Se}_4$ .**

temperature [69]. The value of ZT for  $\text{HgY}_2\text{S}_4$  is 0.708 and 0.742 for  $\text{HgY}_2\text{Se}_4$  at 200K. When temperature increases, this value reaches 0.758 for  $\text{HgY}_2\text{S}_4$  and 0.765 for  $\text{HgY}_2\text{Se}_4$  at 800K. The higher the ZT values, the more will be the thermal-to-electrical conversion efficiency of the material. These investigations revealed that these compositions are potential materials for high-temperature thermoelectric devices. Comparatively speaking, if a thermoelectric device is fabricated using  $\text{HgY}_2\text{Se}_4$  material it will reveal better thermoelectric performance even in a higher temperature range i.e., up to 800K.

#### 4. Conclusion

In nut shell, here we theoretically investigated the mechanical, electronic band structure and transport properties of  $\text{HgY}_2\text{S}/\text{Se}_4$  by DFT-based software WIEN2k. The ground state stability was probed through formation energy, which manifests that  $\text{HgY}_2\text{S}_4$  (−8.64 eV) was more stable than  $\text{HgY}_2\text{Se}_4$  (−6.88 eV). The incorporation of Se at the S site expanded the unit cell and caused an increase in lattice constant from 11.12 to 11.62 Å. In this scenario, the bulk modulus was decreased from 76.19 to 63.65 GPa. Since calculated values of  $B/G > 1.75$  and  $\nu > 0.26$ , uncovered the ductile nature of these compositions. The small direct bandgap semiconducting nature at  $\Gamma$ - $\Gamma$  symmetry point is exposed by electronic bandgap analysis. In addition, keeping S at Se place caused the decrease in bandgap value from 1.2 to 0.6 eV, which was suitable for optical devices operating in the infrared region. The computation of thermoelectric properties has been carried out using the BoltzTrap semiclassical approach. The calculation of transport parameters includes the electronic conductivity, thermal conductivity, Seebeck coefficient, power factor, and figure of merit. Based on all these facts, we conclude that the understudy

spinel chalcogenides are a potential candidate for thermoelectric applications.

#### Data availability statement

All data presented in this work can be provided on request.

#### Declaration of competing interest

The authors declare that they have no known competing financial interests or personal relationships that could have appeared to influence the work reported in this paper.

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