



Dynamic pathways for the bioconvection in thermally activated rotating system

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Received: 30 March 2022 / Revised: 10 June 2022 / Accepted: 15 June 2022

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Abstract

This article communicates Cattaneo-Christov heat and mass flux theory in the flow of Maxwell nanofluid and gyrotactic microorganisms towards a spinning disk with the effects of Hall current, and Darcy porous medium. Entropy generation is also taken into account in the presence of Brownian motion, thermophoresis, thermal radiation, heat source/sink, and chemical reaction. Non-dimensional equations are derived for the velocities, temperature, nanoparticles concentration, and motile gyrotactic microorganisms concentration with the help of similarity transformations. Homotopy analysis method (HAM) is used to obtain the solution. Graphs and a table are plotted to show the effects of physical parameters. Numerical computations are given to analyze the values of skin friction coefficients which present a nice agreement with the published results.

Keywords Cattaneo-Christov heat and mass flux theory · Maxwell nanofluid · Gyrotactic microorganisms · Entropy generation · Homotopy analysis method

Nomenclature

m	Hall parameter	Pr	Prandtl number
(u, v, w)	Velocity components	E'''_{gen}	Entropy generation
(r, ϑ, z)	Cylindrical coordinates	E'''_0	Characteristic entropy generation
a	Stretching rate	N_G	Non-dimensional entropy generation rate
Sc	Schmidt number	Le	Lewis number
k_r	Dimensional chemical reaction parameter	Pe	Peclet number
k_c	Non-dimensional chemical reaction parameter	Nb	Brownian motion parameter
M	Magnetic field parameter	Nt	Thermophoresis parameter
		R	Ideal gas constant
		Br	Brinkman number
		$C_i, i = 1, 2, 3, \dots, 11$	Arbitrary constants
		Re	Reynolds number

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T	Temperature
C	Nanoparticles concentration
N	Motile microorganisms concentration
P	Pressure
c_p	Specific heat at constant pressure
D_B	Diffusivity
B	Chemical species
$f'(\zeta)$	Dimensionless radial velocity
$g(\zeta)$	Dimensionless azimuthal velocity
$h(\zeta)$	Dimensionless axial velocity
B_0	Applied magnetic field strength
L	Linear operator

Greek symbols

Ω	Angular velocity
σ	Electrical conductivity
ψ	Stream function
ζ	Similarity variable
$\phi(\zeta)$	Dimensionless nanoparticles concentration
$\theta(\zeta)$	Dimensionless temperature
$\chi(\zeta)$	Dimensionless gyrotactic microorganisms concentration
λ_1	Relaxation time parameter
γ_1	Thermal relaxation time parameter
γ_2	Mass relaxation time parameter
γ_3	Non-dimensional thermal relaxation time parameter
γ_4	Non-dimensional solutal time relaxation parameter
γ_5	Gyrotactic microorganisms concentration difference parameter
γ_6	Temperature difference parameter
γ_7	Diffusion constant parameter
ν	Kinematic viscosity
μ	Dynamic viscosity
ρ	Density
α	Thermal diffusivity
Ω_1	Dimensionless stretching parameter

Subscripts

f	Base fluid
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Superscripts

'	Differentiation with respect to ζ
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1 Introduction

Non-Newtonian fluids have applications in concentrated products, including fabrics carbon, glass, and paints. Researchers have keen interest in the investigations of non-Newtonian fluids. Wang et al. [1] used the ramped conditions and three different types of definitions of

fractional differential operators, namely Caputo, Caputo-Fabrizio, and Atangana-Baleanu to analyze the Oldroyd-B nanofluid in a porous medium. Farooq et al. [2] explored the characteristics of bioconvection in Carreau nanofluid flow with Cattaneo-Christov heat and mass transfer in the presence of thermal radiation, activation energy, Brownian motion, and thermophoresis. Alzahrani et al. [3] investigated the convective micropolar fluid through a porous nonlinear stretching sheet with variable thickness using the numerical second-order finite difference method, namely the Keller-box scheme. Ahmad et al. [4] examined the steady bio-convective hybridized micropolar nanofluid flow with the stratification conditions above a vertical exponentially stretching surface to analyze the heat and mass transfer rate, activation energy using Cattaneo-Christov heat flux theory. The non-Newtonian behaviors with other characteristics of fluids and heat transfer can be seen in the references [5–14].

Nanofluid is a subclass of liquid consisting of nanometer-sized (1–100 nm) materials, known as nanoparticles. Nanofluid describes the suspended nano-materials in the conventional liquid to enhance the thermal conduction. The nanoparticle use has drawn the attention of investigators in present time due to their remarkable role in the fields of mechanical science, electronics, food processing, biology, medicine, and biomaterials. The targeted drug delivery and cancer treatment hypotheses, fermentation process, nanotechnologies, and medication are also dependent on the applications of the nanofluids. It is a fact that in a refrigerator and other devices, fluid passages through the microfluidics are completely dependent on the heat transporting particles. So nanofluid dynamics is the main term to be understood in all disciplines concerned with nanoparticles suspensions. Khan et al. [15] focused on the boundary layer flow of hybrid nanofluid with gyrotactic microorganisms and porous medium across the rotating disk by considering the double diffusion theory, Joule heating effect, chemical reaction, and multiple slip boundary conditions. Ahmad et al. [16] discussed the unsteady three-dimensional bio-convective Maxwell nanofluid flow through an exponentially stretched sheet with the convective boundary conditions by using Matlab's `bvp4c` function. In another study, Ahmad et al. [17] examined the two different nanoparticles that offer better heat transfer performance and thermophysical properties than traditional heat transfer fluids (water, oil, and refrigerant-134A) and nanofluid with single nanoparticles. Similarly, investigating the nanofluid, Ahmad et al. [18] worked on the unsteady bio-convective hybridized micropolar nanofluid flow over vertical exponentially stretching surface under the stratification conditions with the involvement of Cattaneo-Christov heat flux and activation energy. Khan et al. [38] analyzed bioconvection nanofluid flow and

entropy generation in rotating system with the application of Homotopy Analysis Method. Some relevant materials can be studied in the references [39–43].

Bioconvection is an interesting process of materials associated with the swimming of microorganisms. The phenomenon of bioconvection performs a key role in the production of biochemical oils and meteorological applications including heated spring colonized through motile gyrotactic microorganisms termed as thermopile. Ramesh et al. [44] considered the applications of nano-particles and bioconvection in the flow of Maxwell nanofluid along with swimming of gyrotactic microorganisms configured by a Riga surface in the presence of thermal radiation, activation energy and the convective-Nield boundary conditions. Liu [45] scrutinized the MHD Maxwell nanofluid flow over an extended cylinder with nonlinear thermal radiation amalgamated with chemical reaction in a Darcy-Forchheimer porous medium and gyrotactic microorganisms. Ramzan et al. [46] discussed the new mathematical model of three-dimensional tangent hyperbolic nanofluid flow with Hall current and ion slip in the presence of modified Fourier law with Newtonian heating, Arrhenius activation energy, chemical reaction, and gyrotactic microorganisms. Song et al. [47] reported the stagnation point flow of radiative micropolar nanofluid over an off-centered rotating disk with applications of motile microorganisms incorporating the novel dynamic of thermal radiation and activation energy. Investigating bioconvection, Song et al. [48] communicated the thermal assessment of Sutterby nanofluid containing the gyrotactic microorganisms with solutal and Marangoni boundaries, melting phenomenon, and thermal conductivity. Abdal et al. [49] analyzed the flow of Maxwell nanofluid over an extending sheet with bioconvection of micron size self-motivated organisms, radiative heat flux, and two temperature boundary conditions, namely, prescribed surface temperature (PST) and prescribed heat flux (PHF) under the influence of magnetic field and porosity. Chu et al. [50] highlighted the newly developed concept of Rosseland approximation and gyrotactic microorganisms in steady, two-dimensional, incompressible flow of Walter's-B nanofluid over a stretched surface. The other studies related to bioconvection can be seen in [51–60].

All the above cited papers have remarkable investigations about different fluid flows with heat and mass transfer describing the various characteristics of fluids and effects of parameters on these flows. However, still there is a gap to investigate the Hall current and Darcy Forchheimer porous medium effects with entropy generation as well as gyrotactic microorganisms in rotating system for non-Newtonian nanofluid. The current study presents the investigation about Hall current effect, porosity, and entropy generation in the Maxwell nanofluid flow past a stretchable rotating disk in the presence of Cattaneo-Christov heat and mass flux theory.

The problem is mathematically modeled and then solved through an analytical technique [61, 62].

2 Methods

2.1 Basic equations

The steady axisymmetric swirling flow of incompressible Maxwell nanofluid containing nanoparticles and motile gyrotactic microorganisms in the presence of Hall current effect and entropy generation is investigated. The Cattaneo-Christov theory is applied for the heat and mass transfer by assuming the flow in the upper plane $z \geq 0$, with uniform angular velocity, stretching rate, constant temperature, nanoparticles concentration, and motile gyrotactic microorganisms concentration as Ω , a , T_w , C_w , and N_w respectively while at the free stream, the temperature, nanoparticles concentration, and motile gyrotactic microorganisms concentration are T_∞ , C_∞ , and N_∞ respectively. Magnetic field is applied with the power B_0 along with the z -direction (please see Fig. 1). For the life of microorganisms, water is taken as the base fluid.

Following the above assumptions, governing equations for the conservation of total mass, flow, heat transfer, nanoparticles concentration, and motile microorganisms are presented as in [23, 24, 38].

$$\frac{\partial u}{\partial r} + \frac{u}{r} + \frac{\partial w}{\partial z} = 0, \quad (1)$$

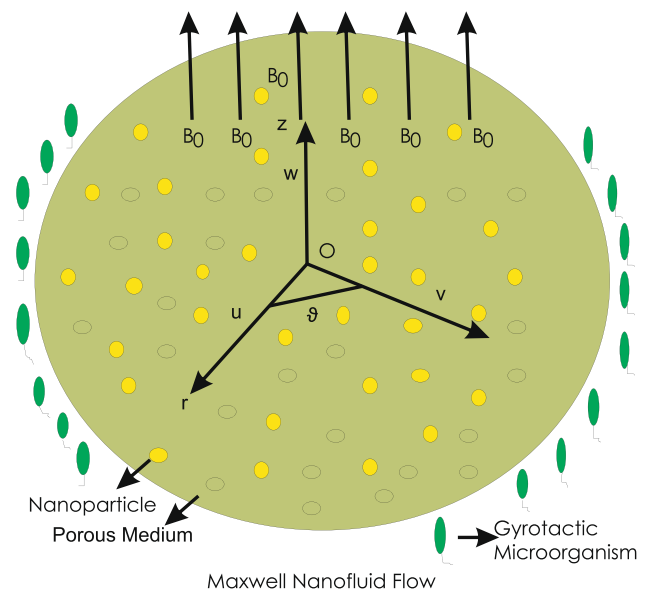


Fig. 1 Physical configuration of the problem

$$\rho_f \left(u \frac{\partial u}{\partial r} + w \frac{\partial u}{\partial z} - \frac{v^2}{r} \right) = \mu_f \frac{\partial^2 u}{\partial z^2} - \lambda_1 \left[u^2 \frac{\partial^2 u}{\partial r^2} + w^2 \frac{\partial^2 u}{\partial z^2} + 2uw \frac{\partial^2 u}{\partial r \partial z} - \frac{2uv}{r} \frac{\partial v}{\partial r} - \frac{2vw}{r} \frac{\partial v}{\partial z} + \frac{uv^2}{r^2} + \frac{v^2}{r} \frac{\partial u}{\partial r} \right] - \frac{\sigma_f B_0^2 (u - mv + w \lambda_1 \frac{\partial u}{\partial z})}{1 + m^2} - \frac{\mu_f}{k} u - k^F u^2, \quad (2)$$

$$\rho_f \left(u \frac{\partial v}{\partial r} + w \frac{\partial v}{\partial z} + \frac{1}{r} uv \right) = \mu_f \frac{\partial^2 v}{\partial z^2} - \lambda_1 \left[u^2 \frac{\partial^2 v}{\partial r^2} + w^2 \frac{\partial^2 v}{\partial z^2} + 2uw \frac{\partial^2 v}{\partial r \partial z} + \frac{2uv}{r} \frac{\partial u}{\partial r} + \frac{2vw}{r} \frac{\partial u}{\partial z} - \frac{2u^2}{r^2} - \frac{v^3}{r^2} + \frac{v^2}{r^2} \frac{\partial v}{\partial r} \right] - \frac{\sigma_f B_0^2 (v + mu + w \lambda_1 \frac{\partial v}{\partial z})}{1 + m^2} - \frac{\mu_f}{k} w - k^F w^2, \quad (3)$$

$$\left(u \frac{\partial T}{\partial r} + w \frac{\partial T}{\partial z} \right) = \frac{\alpha}{\rho c_p} \frac{\partial^2 T}{\partial z^2} + 2\gamma_1 \tau \frac{D_T}{T_\infty} \left[u \frac{\partial T}{\partial z} \frac{\partial^2 T}{\partial r \partial z} + w \frac{\partial T}{\partial z} \frac{\partial^2 T}{\partial z^2} \right] + \tau \left[D_B \left(\frac{\partial T}{\partial z} \frac{\partial C}{\partial z} \right) + \frac{D_T}{T_\infty} \left(\frac{\partial T}{\partial z} \right)^2 \right] + \frac{Q_0}{\rho c_p} (T - T_\infty) + \gamma_1 \frac{Q_0}{\rho c_p} \left(u \frac{\partial T}{\partial r} + w \frac{\partial T}{\partial z} \right) + \gamma_1 \tau D_B \left[u \frac{\partial T}{\partial z} \frac{\partial^2 C}{\partial r \partial z} + u \frac{\partial C}{\partial z} \frac{\partial^2 T}{\partial r \partial z} + w \frac{\partial T}{\partial z} \frac{\partial^2 C}{\partial z^2} + w \frac{\partial C}{\partial z} \frac{\partial^2 T}{\partial z^2} \right] + \gamma_1 \left[u \frac{\partial w}{\partial r} \frac{\partial T}{\partial z} + u \frac{\partial u}{\partial r} \frac{\partial T}{\partial r} + w \frac{\partial u}{\partial z} \frac{\partial T}{\partial r} + w \frac{\partial w}{\partial z} \frac{\partial T}{\partial z} + u^2 \frac{\partial^2 T}{\partial r^2} + w^2 \frac{\partial^2 T}{\partial z^2} + 2uw \frac{\partial^2 T}{\partial r \partial z} \right], \quad (4)$$

$$u \frac{\partial C}{\partial r} + w \frac{\partial C}{\partial z} = D_B \frac{\partial^2 C}{\partial z^2} + \gamma_2 \tau \frac{D_T}{T_\infty} \left[u \frac{\partial^3 T}{\partial r \partial z^2} + w \frac{\partial^3 T}{\partial z^3} \right] + \frac{D_T}{T_\infty} \frac{\partial^2 T}{\partial z^2} - k_r (C - C_\infty) + \gamma_2 k_r \left(u \frac{\partial C}{\partial r} + w \frac{\partial C}{\partial z} \right) - \gamma_2 \left[u \frac{\partial w}{\partial r} \frac{\partial C}{\partial z} + u \frac{\partial u}{\partial r} \frac{\partial C}{\partial r} + w \frac{\partial u}{\partial z} \frac{\partial C}{\partial r} + w \frac{\partial w}{\partial z} \frac{\partial C}{\partial z} + u^2 \frac{\partial^2 C}{\partial r^2} + w^2 \frac{\partial^2 C}{\partial z^2} + 2uw \frac{\partial^2 C}{\partial r \partial z} \right], \quad (5)$$

$$u \frac{\partial N}{\partial r} + w \frac{\partial N}{\partial z} + \frac{bW_c}{(C_w - C_\infty)} \frac{\partial}{\partial z} \left(N \frac{\partial C}{\partial z} \right) = D_n \left[\frac{\partial^2 N}{\partial r^2} + \frac{1}{r} \frac{\partial N}{\partial r} + \frac{\partial^2 N}{\partial z^2} \right], \quad (6)$$

using the boundary conditions

$$u = ra, \quad v = r\Omega, \quad w = 0, \quad T = T_w, \quad C = C_w, \quad N = N_w \quad \text{at } z = 0, \quad (7)$$

$$u \rightarrow 0, \quad v \rightarrow 0, \quad T \rightarrow T_\infty, \quad C \rightarrow C_\infty, \quad N \rightarrow N_\infty \quad \text{as } z \rightarrow \infty, \quad (8)$$

where $u(r, \vartheta, z)$, $v(r, \vartheta, z)$ and $w(r, \vartheta, z)$ are the velocity components, P is the pressure, $B = (0, 0, B_0)$ is the magnetic induction, and m is the Hall parameter [38]. λ_1 is the relaxation time parameter. γ_1 and γ_2 are the thermal relaxation time and mass relaxation time parameters respectively. α is the thermal diffusivity, k is the permeability of porous medium, and $k^F = \frac{C_b}{\sqrt{k}}$ is the Forchheimer resistance factor in which C_b is the drag coefficient. D_B and D_T are the Brownian and thermophoretic diffusion coefficients, ρ_f , μ_f , σ_f and $(c_p)_f$ are the density, effective dynamic viscosity, electrical conductivity, and heat capacitance of the nanofluid respectively. $\nu_f = \frac{\mu_f}{\rho_f}$ is the kinematic viscosity, $\tau = \frac{(\rho c)_p}{(\rho c)_f}$ is the ratio of heat capacity of nanoparticles to heat capacity of nanofluid, Q_0 is the heat

generation source parameter, k_r is the chemical reaction rate, b is the chemotaxis constant, W_c is the maximum cell swimming speed, and D_n is the diffusivity of microorganisms. at $z = 0$,

In order to make the above governing Eqs. (1–9) dimensionless, the following transformations are introduced

$$u = r\Omega f(\zeta), \quad v = r\Omega g(\zeta), \quad w = (\Omega \nu_f)^{\frac{1}{2}} h(\zeta), \quad \theta(\zeta) = \frac{T - T_\infty}{T_w - T_\infty}, \quad \phi(\zeta) = \frac{C - C_\infty}{C_w - C_\infty}, \quad \chi(\zeta) = \frac{N - N_\infty}{N_w - N_\infty}, \quad \zeta = \left[\frac{\Omega}{\nu_f} \right]^{\frac{1}{2}} z. \quad (9)$$

Substituting the values from Eqs. (9) in (1–8), the following eight Eqs. (10–17) are obtained.

$$2f + h' = 0, \quad (10)$$

$$f^2 - g^2 + f'h - f'' + \beta_1 \left(h^2 f'' + 2ff'h - 2gg'h \right) - \frac{M(f' - mg - 2\beta_1 hf')}{1 + m^2} - \lambda_3 f - \lambda_4 f^2 = 0, \quad (11)$$

$$2fg + g'h - g'' + \beta_1 (g''h^2 + 2(fg' + f'g)h) - \frac{M(mf' + g - 2\beta_1hg')}{1 + m^2} - \lambda_3g - \lambda_4g^2 = 0, \quad (12)$$

$$\frac{1}{Pr}\theta'' - h\theta' + Nb(\theta'\phi' + \gamma_3(h\theta'\phi'' + h\theta''\phi')) + Nt((\theta')^2 + 2\gamma_3h\theta'\theta'') - \gamma_3(h^2\theta'' + hh'\theta')\beta_2\theta + \gamma_3\beta_2h\theta' = 0, \quad (13)$$

$$\phi'' - PrLe\phi' - \gamma_4PrLe(h^2\phi'' + hh'\phi') + \frac{Nt}{Nb}(\theta'' + \gamma_4h\theta''') - PrLe\phi - PrLe\phi h\phi' = 0, \quad (14)$$

$$\chi'' - Lbh\chi' - Pe(\chi'\phi' + \phi''(\gamma_5 + \chi)) = 0, \quad (15)$$

$$f = \Omega_1, g = 1, h = 0, \theta = 1, \phi = 1, \chi = 1 \quad \text{at} \quad \zeta = 0, \quad (16)$$

$$f \rightarrow 0, g \rightarrow 0, h \rightarrow 0, \theta \rightarrow 0, \phi \rightarrow 0, \chi \rightarrow 0 \quad \text{as} \quad \zeta \rightarrow \infty, \quad (17)$$

where (') represents the differentiability through ζ . $\beta_1 = \lambda_1\Omega$ is the non-dimensional Maxwell nanofluid parameter, $\gamma_3 = \gamma_1\Omega$ is the non-dimensional thermal relaxation time parameter, $\gamma_4 = \gamma_2\Omega$ is the non-dimensional solutal relaxation time parameter, $\Omega_1 = \frac{a}{\Omega}$ is the stretching parameter, $\lambda_3 = \frac{\nu_f}{k\Omega}$ is the porosity parameter, $\lambda_4 = \frac{k^F}{\sqrt{k}}$ is the Darcy Forchheimer parameter, $M = \frac{\sigma_f B_0^2}{\rho_f \Omega}$ is the magnetic field parameter, $Pr = \frac{\nu_f}{\alpha}$ denotes the Prandtl number, $\beta_2 = \frac{Q_0}{(\rho c_p)\Omega}$ is the heat generation parameter, $k_c = k_1\Omega$ is the chemical reaction parameter, $Lb = \frac{\nu_f}{D_n}$ is the bioconvection Lewis number, $Le = \frac{\nu_f}{D_{B_\infty}}$ is the Lewis number, $Pe = \frac{bW_c}{D_n}$ is the Peclet number, $\gamma_5 = \frac{N_w - N_\infty}{N_w - N_\infty}$ is microorganisms concentration difference parameter, $Nb = \frac{\tau D_B(C_w - C_\infty)}{\nu_f T_\infty}$ is the Brownian motion parameter, and $Nt = \frac{\tau D_T(T_w - T_\infty)}{\nu_f T_\infty}$ is the thermophoresis parameter.

The quantities of engineering interest, such as coefficient of skin friction C_F , local Nusselt Nu_r , Sherwood number Sh_r , and local motile density number Nn_r , are defined as

$$C_F = \frac{\tau|_{\zeta=0}}{\rho_f(r\Omega)^2}, \quad (18)$$

where

$$\tau = \sqrt{(\tau_r)^2 + (\tau_\theta)^2}, \quad (19)$$

denotes the sum of shear stresses τ_r and τ_θ along radial and transvers directions.

$$Nu_r = \frac{-rq_1}{\alpha(T_w - T_\infty)}, \quad Sh_r = \frac{-rq_2}{D_B(C_w - C_\infty)}, \quad Nn_r = \frac{-rq_3}{D_m(N_w - N_\infty)}, \quad (20)$$

where q_1, q_2 and q_3 are the heat, mass, and motile micro-organism fluxes at the surface of the rotating disk and are defined as

$$q_1 = -\alpha T_z|_{\zeta=0}, \quad q_2 = -D_B C_z|_{\zeta=0}, \quad q_3 = D_m N_z|_{\zeta=0}. \quad (21)$$

Using the information of Eqs. (9) and (18) proceeds to

$$C_F = Re_r^{-\frac{1}{2}} \left[(f'(0))^2 + (g'(0))^2 \right], \quad (22)$$

where $Re_r = \frac{r^2\Omega}{\nu_f}$ is the Reynolds number.

Similarly by putting values from Eqs. (9) in (20), it is obtained that

$$Nu_r = -Re_r^{0.5}\theta'(0), \quad Sh_r = -Re_r^{0.5}\phi'(0), \quad Nn_r = -Re_r^{0.5}\chi'(0). \quad (23)$$

3 Entropy generation

Through fluid mechanics, in system and its surroundings, entropy generation measures the rate of disorderness. It is related with the friction of particles and kinetic and potential energies. The entropy generation has the equation [51]

$$E_{gen}''' = \left(\frac{\partial T}{\partial z} \right)^2 \frac{\alpha}{T_\infty^2} + \frac{1}{T_\infty} \left[\left(\frac{\partial u}{\partial z} \right)^2 + \left(\frac{\partial v}{\partial z} \right)^2 \right] + \left(\frac{\partial C}{\partial z} \right)^2 \frac{RD}{C_\infty} + \left[\frac{\partial C}{\partial z} \frac{\partial T}{\partial z} \right] \frac{RD}{T_\infty} + \left(\frac{\partial N}{\partial z} \right)^2 \frac{RD}{N_\infty} + \left[\frac{\partial N}{\partial z} \frac{\partial C}{\partial z} \right] \frac{RD}{T_\infty} + \frac{\sigma_f B_0^2 (u^2 + v^2)}{1 + m^2}, \quad (24)$$

where R is the ideal gas constant while D is the diffusivity.

Table 1 Numerical information for the comparison of the present work

Profiles	Ahmed et al. [62]	Present results
$f'(0)$	0.5000776	0.5000775
$g'(0)$	0.6185207	0.6185208
$\theta'(0)$	0.9300460	0.9300461

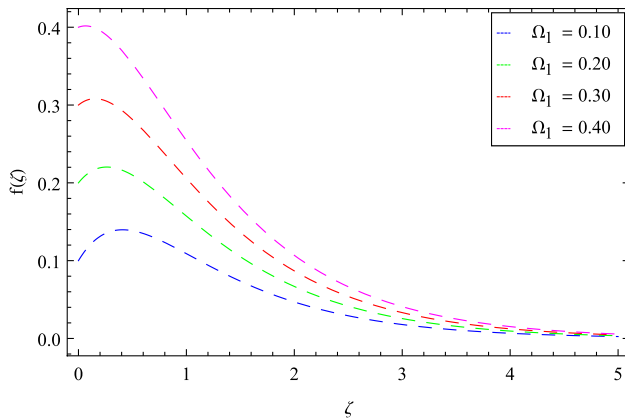


Fig. 2 Behavior of profile against parameter effect

The characteristic entropy generation is

$$E_0''' = \frac{\alpha(T_w - T_\infty)}{\nu_f T_\infty}. \quad (25)$$

The non-dimensional entropy generation rate $N_G(\zeta)$ is obtained with the help of Eq. (9) as

$$N_G(\zeta) = \frac{E_{gen}'''}{E_0'''} \quad (26)$$

So

$$\begin{aligned} N_G(\zeta) = & \gamma_6(\theta')^2 + Br \left[f'^2 + g'^2 \right] \\ & + \gamma_7\theta'\phi' + \gamma_7\frac{\gamma_8}{\gamma_6}(\phi')^2 + \gamma_9\theta'\chi' \\ & + \gamma_9\frac{\gamma_5}{\gamma_6}(\chi')^2 + \frac{BrM}{1+m^2} \left[f^2 + g^2 \right], \end{aligned} \quad (27)$$

where $\gamma_6 = \frac{T_w - T_\infty}{T_\infty}$ is the temperature difference parameter, $Br = \frac{\mu_f \Omega^2}{\alpha(T_w - T_\infty)}$ is the Brinkman number, $\gamma_7 = \frac{RD(C_w - C_\infty)}{\alpha}$ is the diffusion constant parameter due to nanoparticles

concentration, $\gamma_8 = \frac{C_w - C_\infty}{C_\infty}$ is the nanoparticles concentration difference parameter, and $\gamma_9 = \frac{RD(N_w - N_\infty)}{\alpha}$ is the diffusion constant parameter due to microorganisms concentration.

4 Solution methodology

Applying a homotopy analysis method (HAM) [61], the initial approximations and auxiliary linear operators are

$$\begin{aligned} h_0(\zeta) = 0, f_0(\zeta) = \Omega \exp(-\zeta), g_0(\zeta) = \exp(-\zeta), \\ \theta_0(\zeta) = \exp(-\zeta), \phi_0(\zeta) = \exp(-\zeta), \chi_0(\zeta) = \exp(-\zeta), \end{aligned} \quad (28)$$

$$\begin{aligned} L_h = h', L_f = f'' - f', L_g = g'' - g', L_\theta = \theta'' - \theta, \\ L_\phi = \phi'' - \phi, L_\chi = \chi'' - \chi \end{aligned} \quad (29)$$

The following properties are satisfied with the linear operators

$$\begin{aligned} L_h [C_1] = 0, L_f [C_2 \exp(\zeta) + C_3 \exp(-\zeta)] = 0, \\ L_g [C_4 \exp(\zeta) + C_5 \exp(-\zeta)] = 0, \\ L_\theta [C_6 \exp(\zeta) + C_7 \exp(-\zeta)] = 0, \\ L_\phi [C_8 \exp(\zeta) + C_9 \exp(-\zeta)] = 0, \\ L_\chi [C_{10} \exp(\zeta) + C_{11} \exp(-\zeta)] = 0, \end{aligned} \quad (30)$$

where $C_i (i = 1-11)$ are the arbitrary constants.

5 Comparison of the present work with the open literature

For the comparison of the present work, the achieved results for $f'(0)$, $g'(0)$ and $\theta'(0)$ are tabulated in Table 1 which shows the close agreement with [62].

6 Results and discussion

The homotopy analysis method (HAM) is used to solve the non-dimensional Eqs. (10–17) and hence Eq. (27). Parameter effects on the different profiles are shown in the relevant graphs.

Figure 2 shows that by increasing the stretching parameter Ω_1 , the radial velocity $f(\zeta)$ is increased. In fact, stretching parameter Ω_1 is the ratio of stretching constant a to the angular velocity Ω . Therefore by the increasing values of stretching parameter Ω_1 , the radial velocity $f(\zeta)$ increases due to disk stretching rate. The effect of Hall parameter m is shown in Fig. 3. The high values of m in Eq. (11) for the term $\frac{M(f' - mg - 2\beta_1 hf')}{1 + m^2}$ show that radial velocity $f(\zeta)$ is decreased. Figures 4 and 5 reveal the effects of porous and Darcy Forchheimer parameters (λ_3, λ_4) on radial velocity $f(\zeta)$. Fluid motion is enhanced by these two parameters.

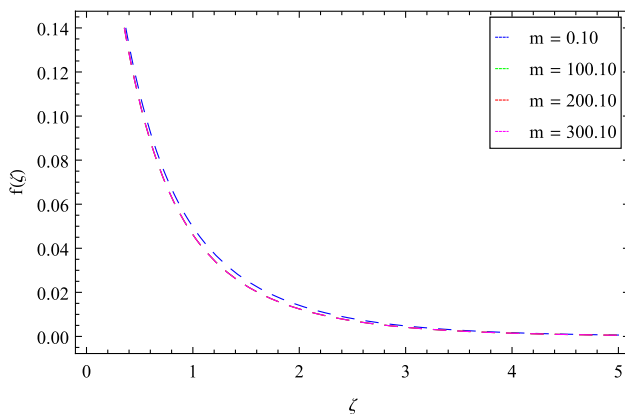


Fig. 3 Behavior of profile against parameter effect

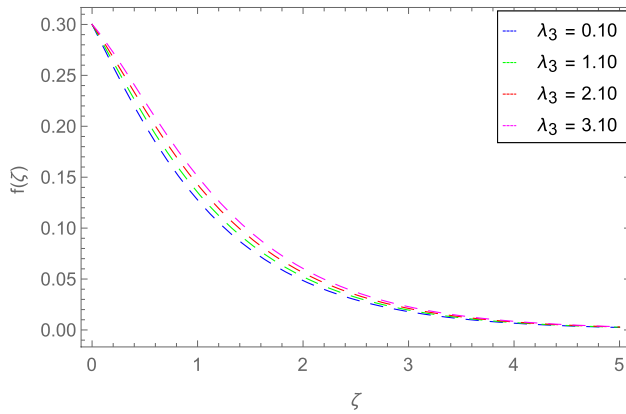


Fig. 4 Behavior of profile against parameter effect

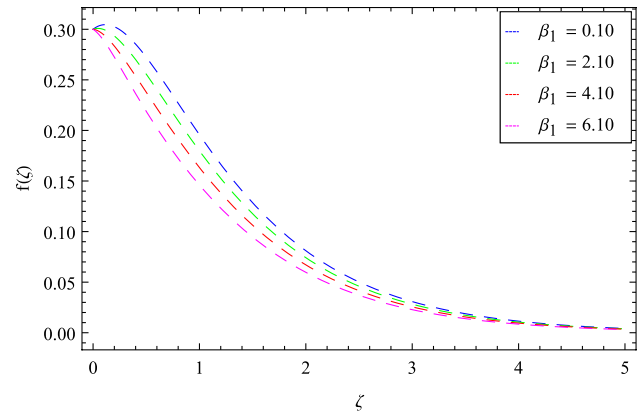


Fig. 7 Behavior of profile against parameter effect

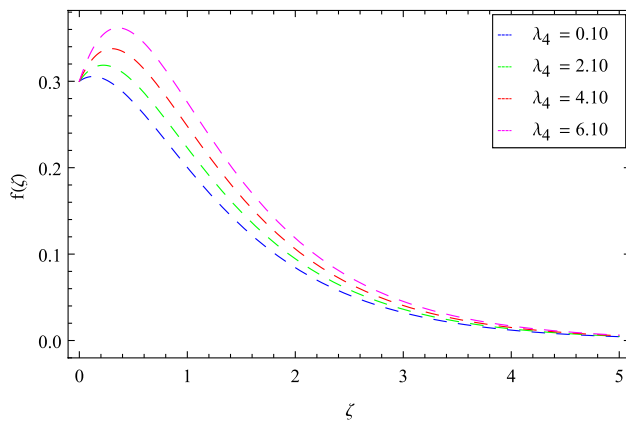


Fig. 5 Behavior of profile against parameter effect

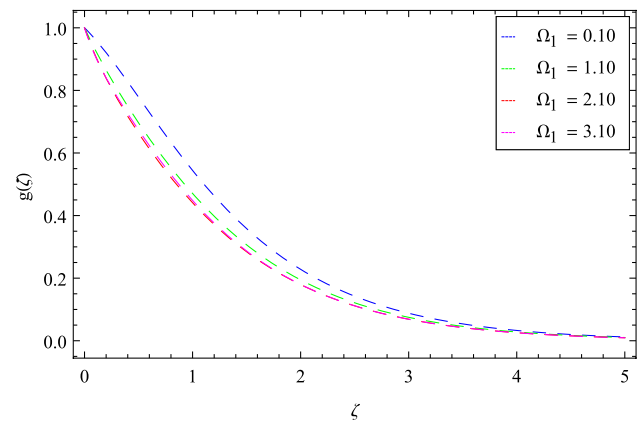


Fig. 8 Behavior of profile against parameter effect

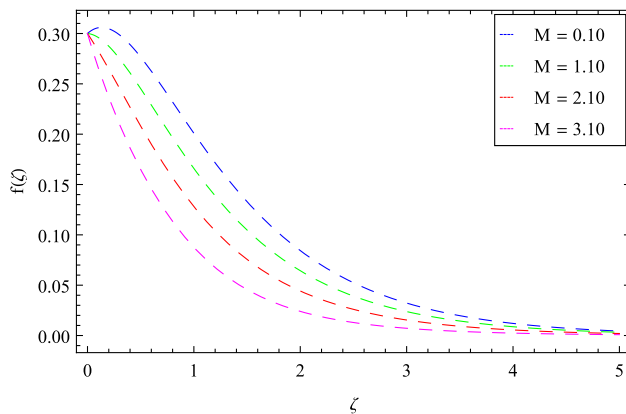


Fig. 6 Behavior of profile against parameter effect

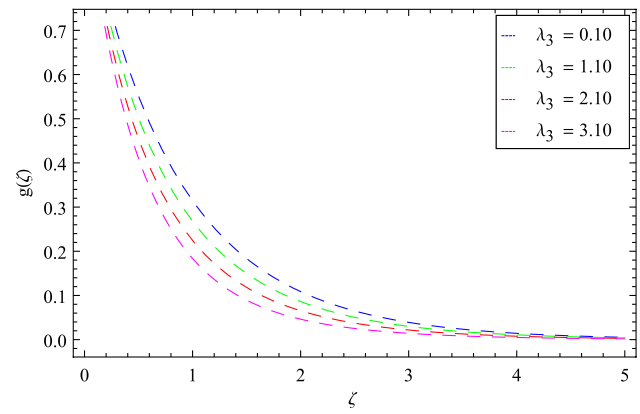


Fig. 9 Behavior of profile against parameter effect

Figure 6 exhibits that the radial velocity $f(\zeta)$ reduces due to magnetic field parameter M . This scenario is on account of Lorentz effect which yields a retarding force along with magnetic field. Figure 7 discloses the role of Maxwell nano-fluid parameter β_1 and the radial velocity $f(\zeta)$. The fluid flow

is reduced because for the greater values of β_1 , the fluid resumes the form of solid which slows down the motion.

Figure 8 is drawn to show that the azimuthal velocity $g(\zeta)$ is reduced with stretching parameter Ω_1 . Similarly in Figs. 9 and 10, it is observed that porous and Darcy Forchheimer parameters have decreasing effects for azimuthal

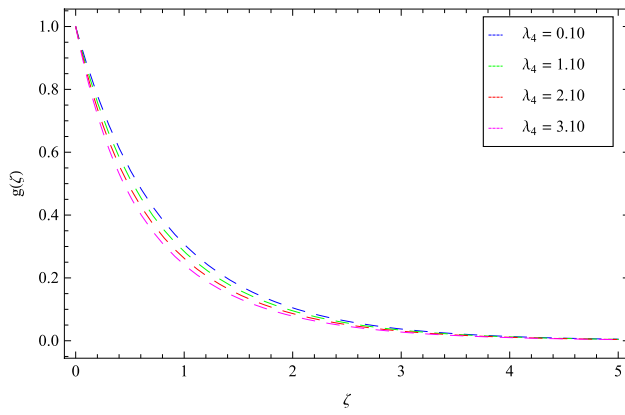


Fig. 10 Behavior of profile against parameter effect

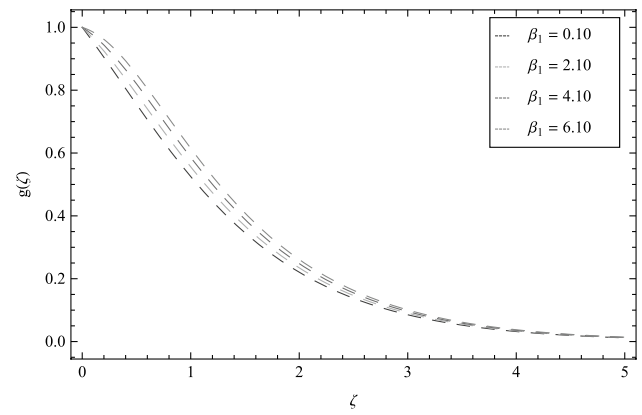


Fig. 13 Behavior of profile against parameter effect

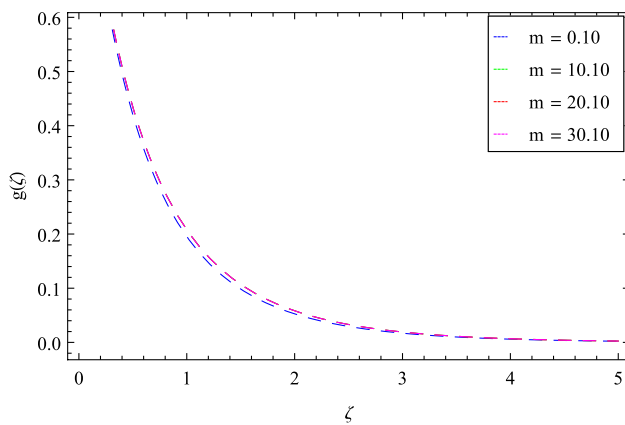


Fig. 11 Behavior of profile against parameter effect

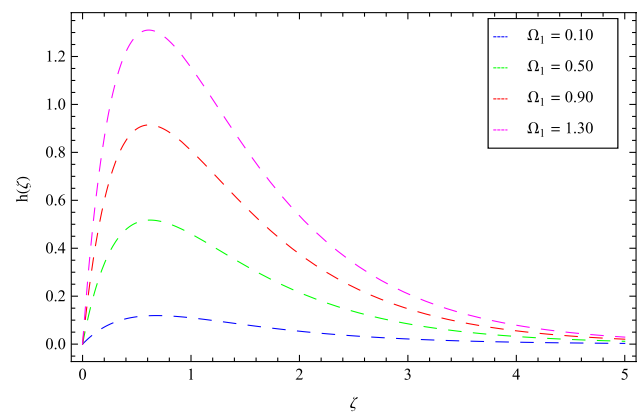


Fig. 14 Behavior of profile against parameter effect

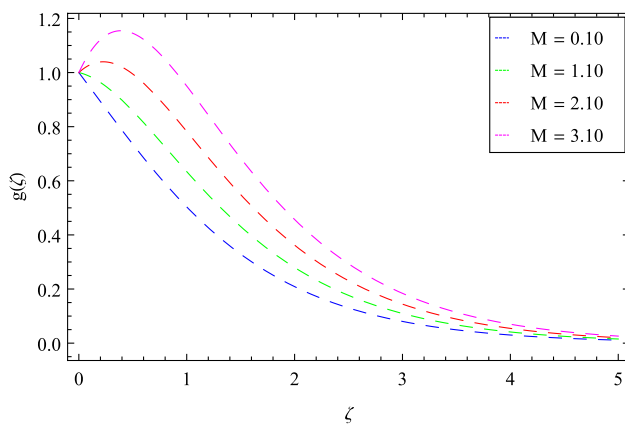


Fig. 12 Behavior of profile against parameter effect

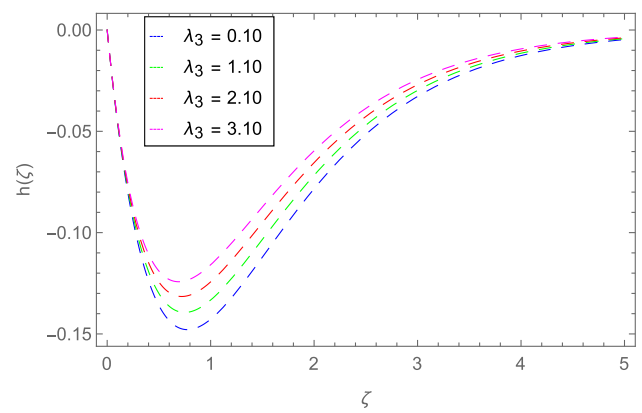


Fig. 15 Behavior of profile against parameter effect

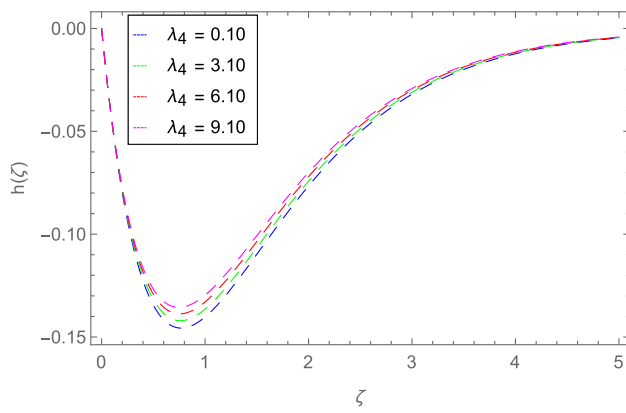


Fig. 16 Behavior of profile against parameter effect

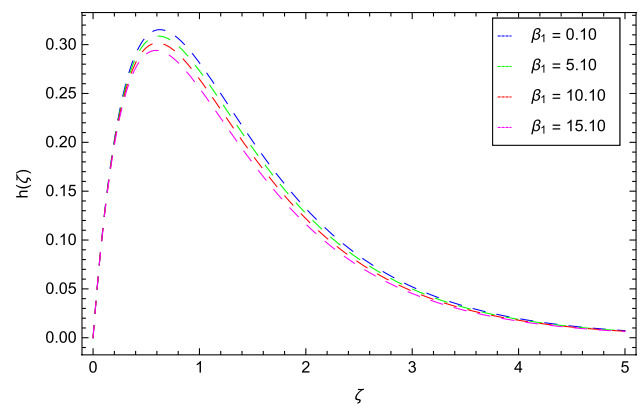


Fig. 19 Behavior of profile against parameter effect

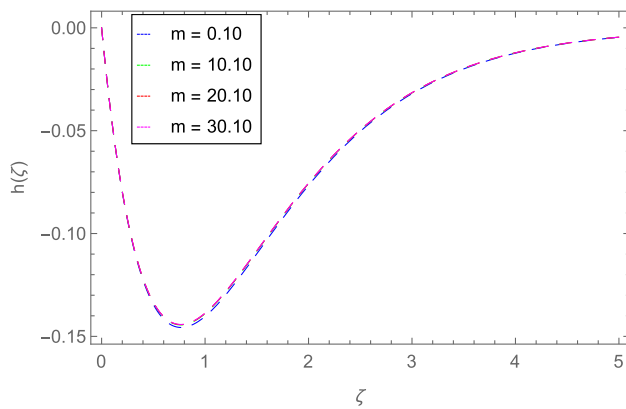


Fig. 17 Behavior of profile against parameter effect

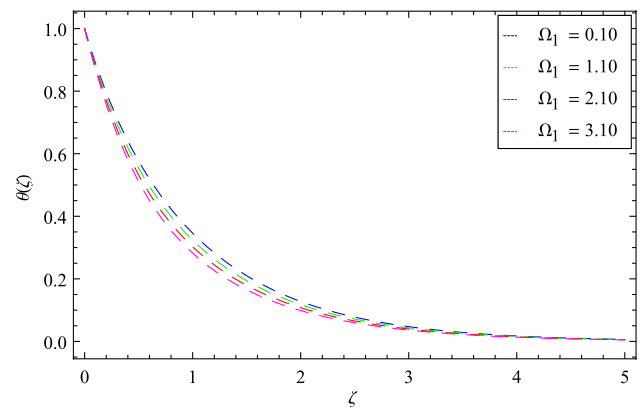


Fig. 20 Behavior of profile against parameter effect

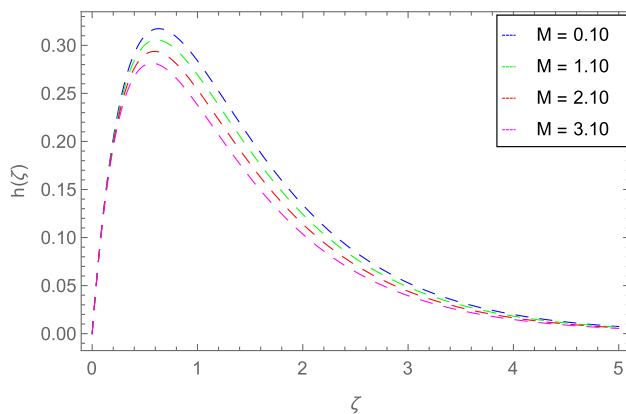


Fig. 18 Behavior of profile against parameter effect

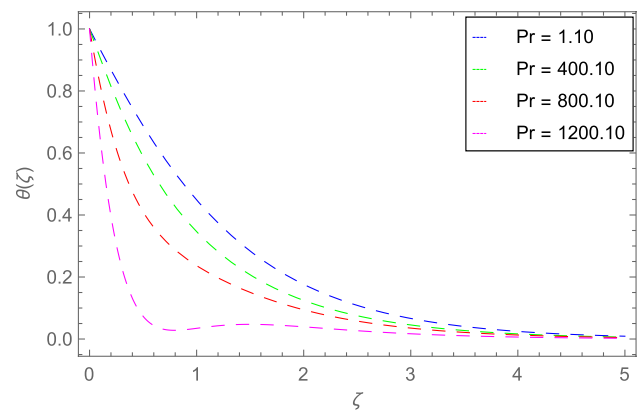


Fig. 21 Behavior of profile against parameter effect

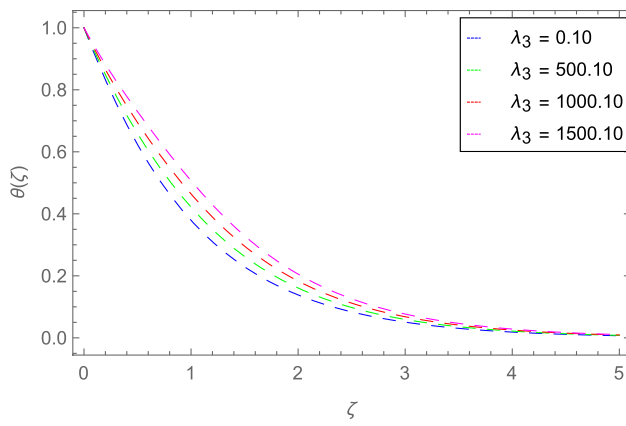


Fig. 22 Behavior of profile against parameter effect

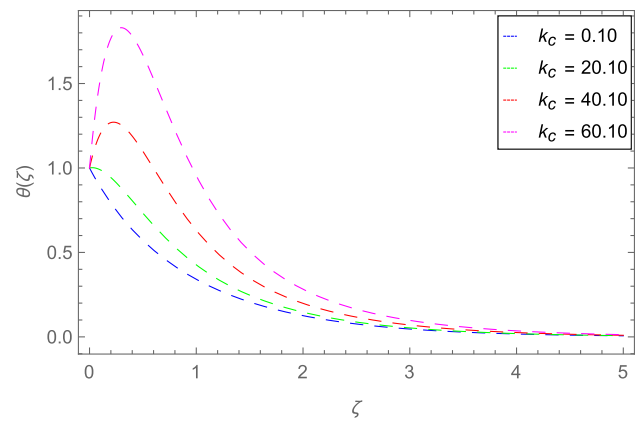


Fig. 25 Behavior of profile against parameter effect

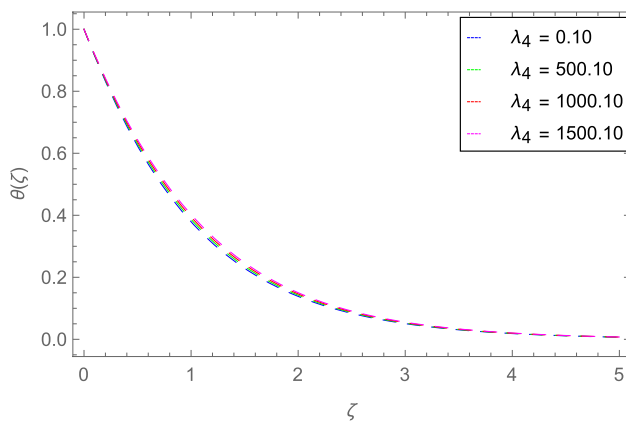


Fig. 23 Behavior of profile against parameter effect

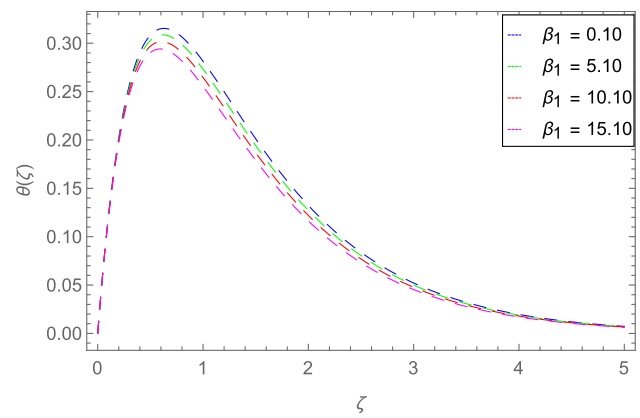


Fig. 26 Behavior of profile against parameter effect

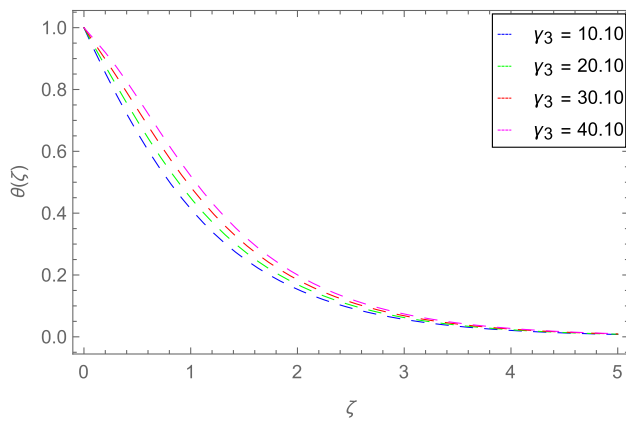


Fig. 24 Behavior of profile against parameter effect

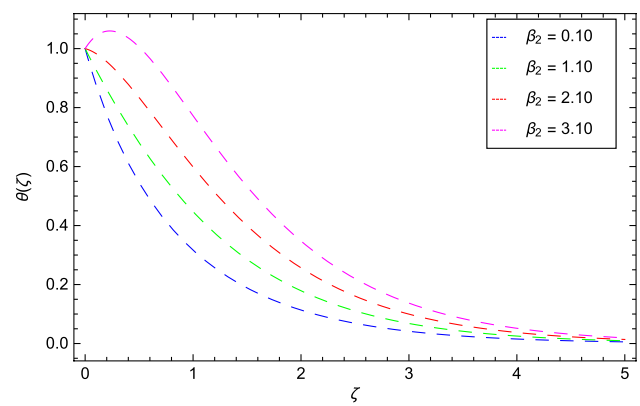


Fig. 27 Behavior of profile against parameter effect

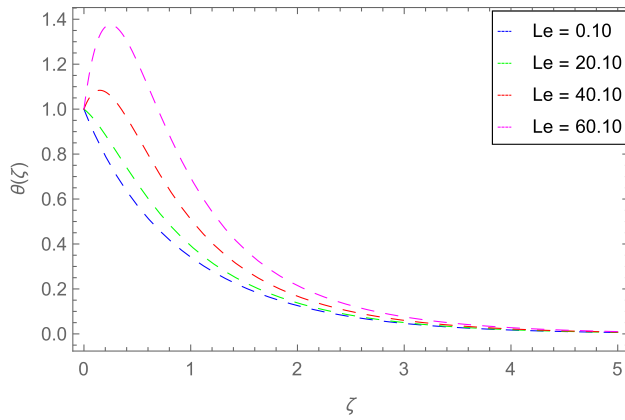


Fig. 28 Behavior of profile against parameter effect

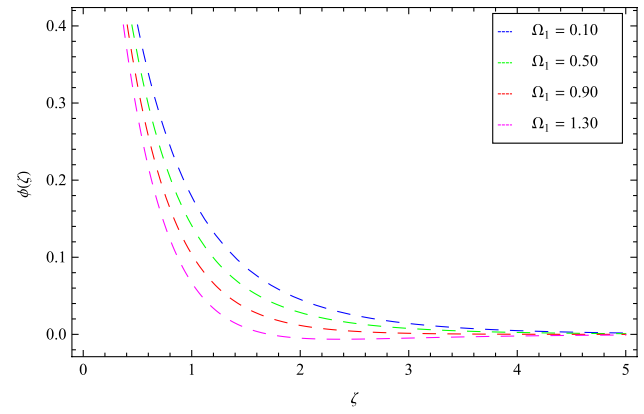


Fig. 31 Behavior of profile against parameter effect

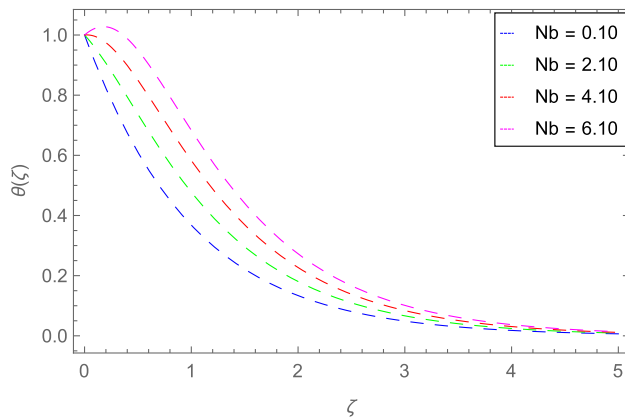


Fig. 29 Behavior of profile against parameter effect

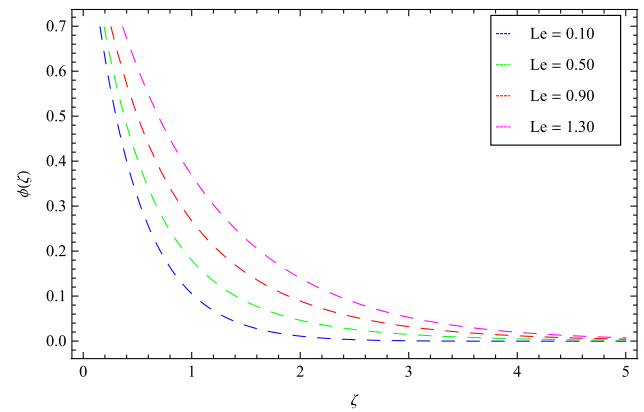


Fig. 32 Behavior of profile against parameter effect

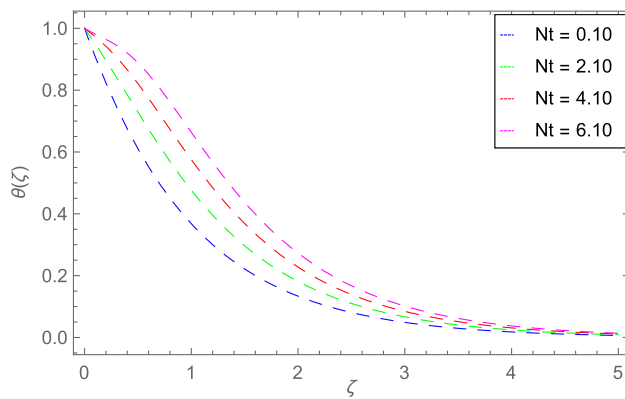


Fig. 30 Behavior of profile against parameter effect

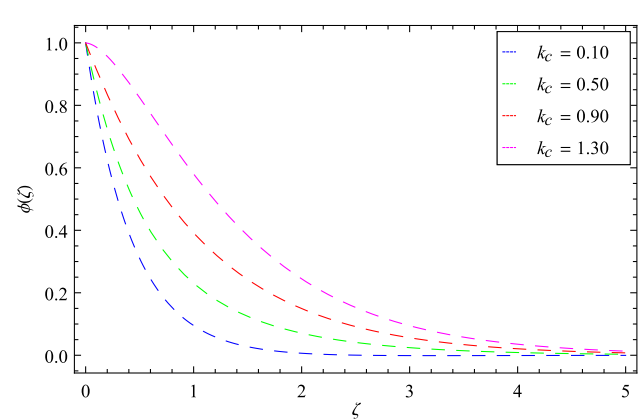


Fig. 33 Behavior of profile against parameter effect

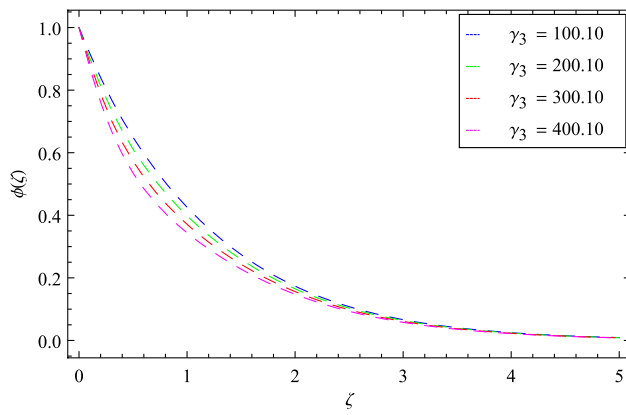


Fig. 34 Behavior of profile against parameter effect

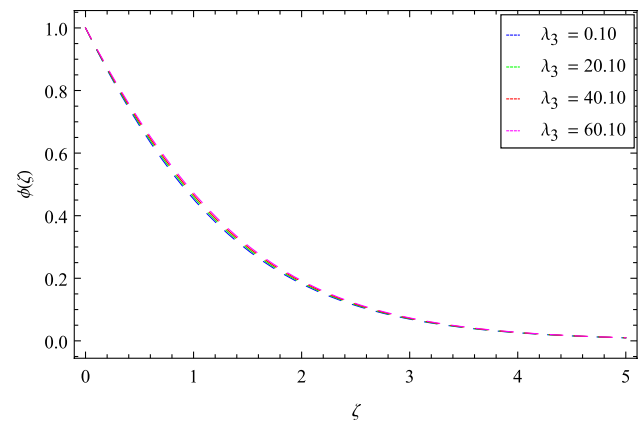


Fig. 37 Behavior of profile against parameter effect

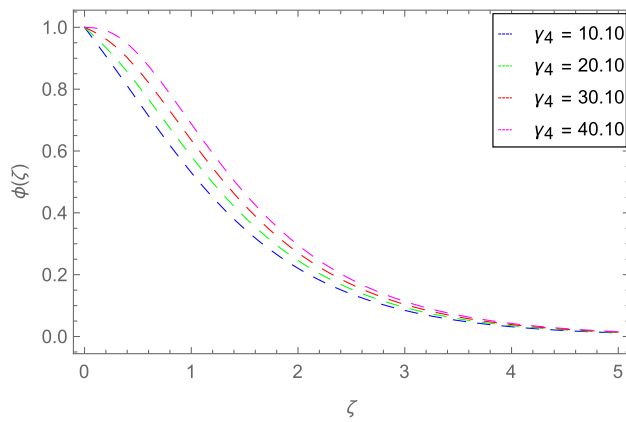


Fig. 35 Behavior of profile against parameter effect

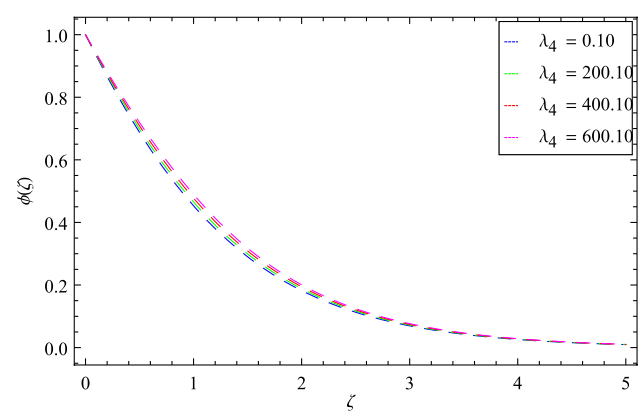


Fig. 38 Behavior of profile against parameter effect

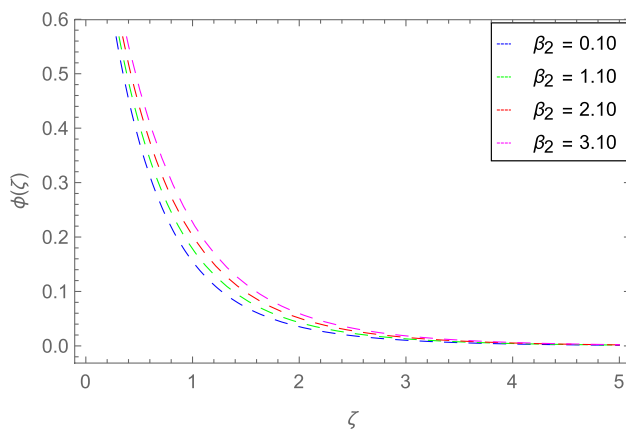


Fig. 36 Behavior of profile against parameter effect

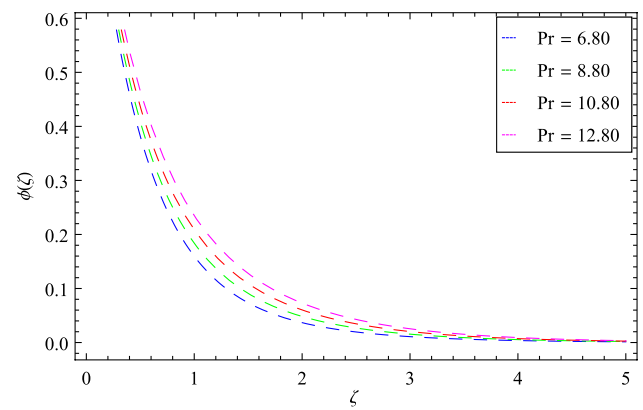


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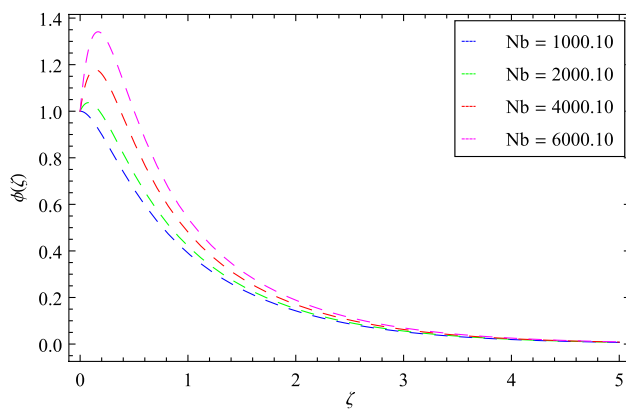


Fig. 40 Behavior of profile against parameter effect

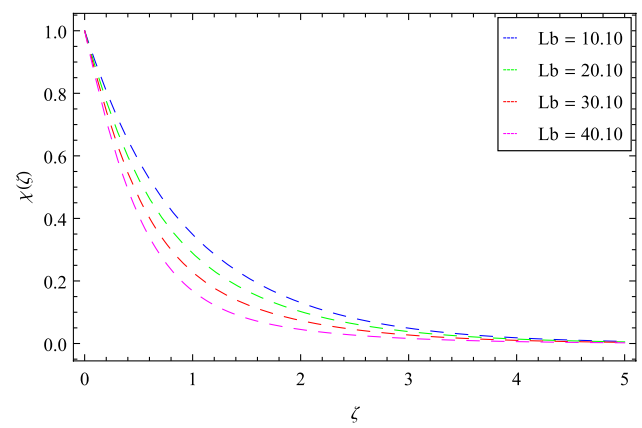


Fig. 43 Behavior of profile against parameter effect

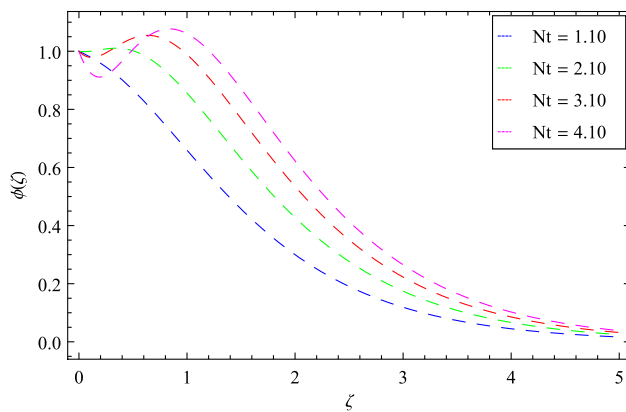


Fig. 41 Behavior of profile against parameter effect

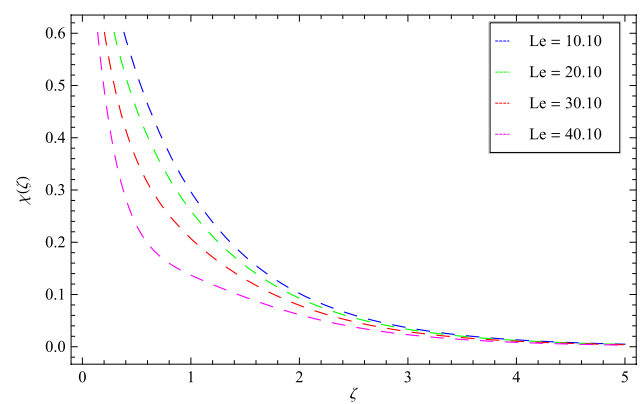


Fig. 44 Behavior of profile against parameter effect

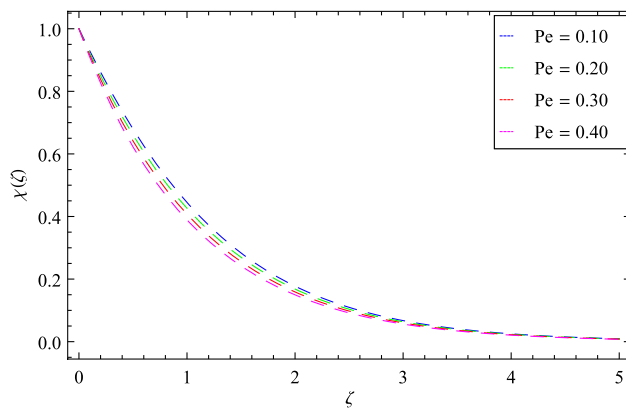


Fig. 42 Behavior of profile against parameter effect

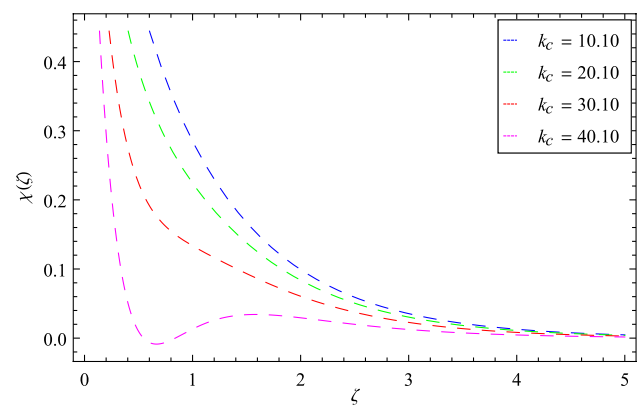


Fig. 45 Behavior of profile against parameter effect

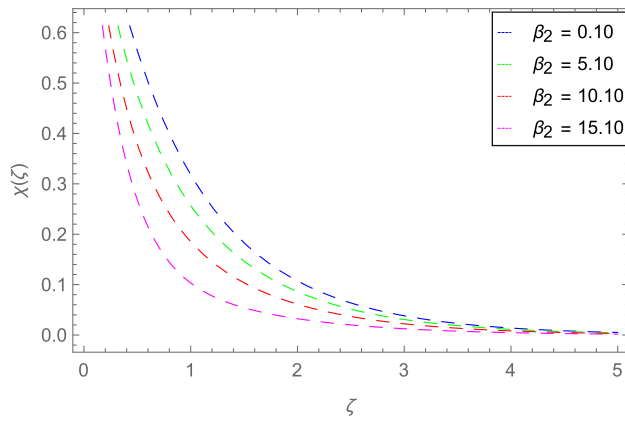


Fig. 46 Behavior of profile against parameter effect

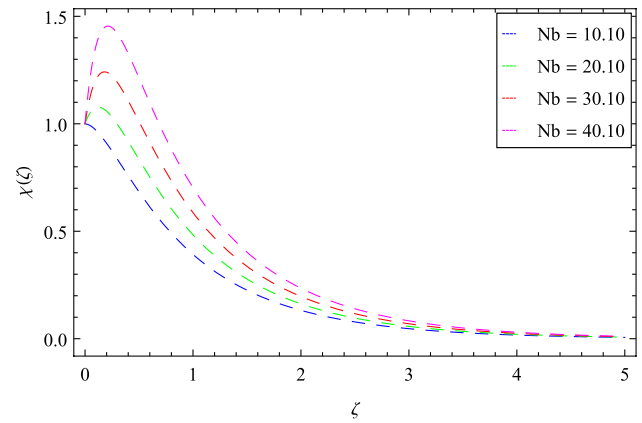


Fig. 49 Behavior of profile against parameter effect

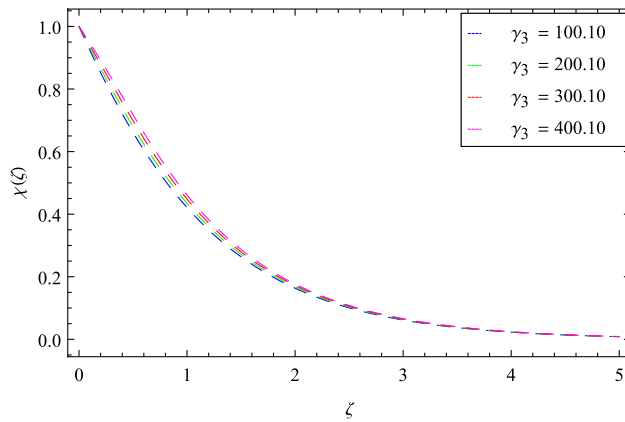


Fig. 47 Behavior of profile against parameter effect

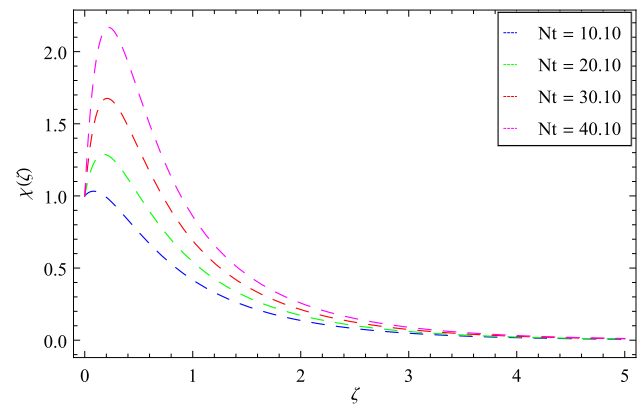


Fig. 50 Behavior of profile against parameter effect

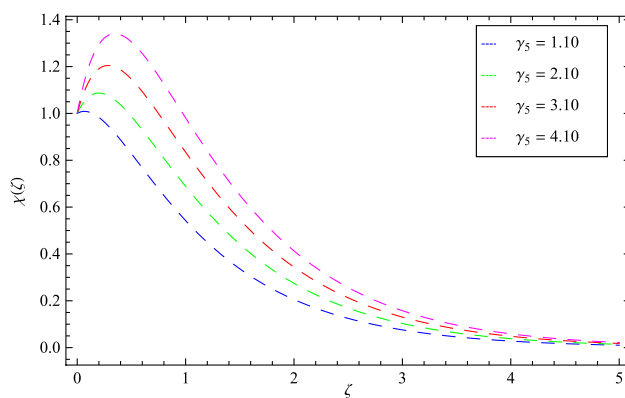


Fig. 48 Behavior of profile against parameter effect

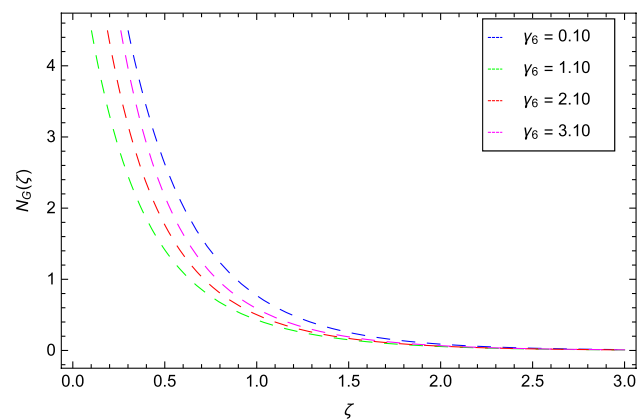
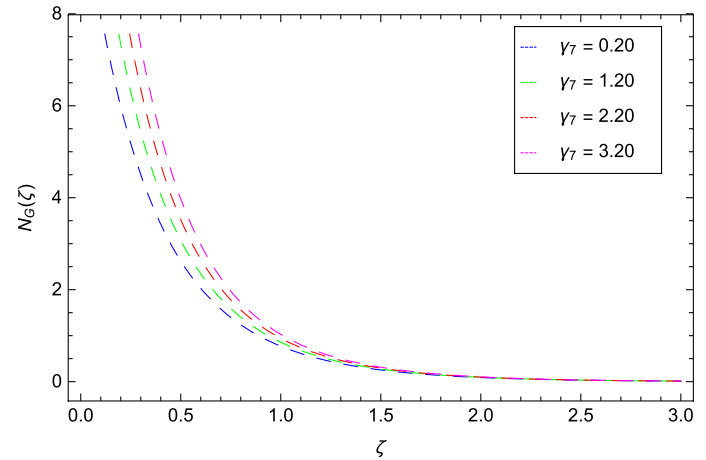


Fig. 51 Behavior of profile against parameter effect

Fig. 52 Behavior of profile against parameter effect



velocity $g(\zeta)$. Fluid motion is reduced by these two parameters (λ_3, λ_4). Figure 11 highlights the effect of Hall parameter m on the azimuthal velocity $g(\zeta)$. It is noted that the velocity component $g(\zeta)$ is made stronger by the given relation $M(mf' + g - 2\beta_1 hg')$ through Eq. (12). The same mathematical relation $\frac{1 + m^2}{M(mf' + g - 2\beta_1 hg')}$ also supports the azimuthal velocity $g(\zeta)$ for the magnetic field parameter M as it is evident from Fig. 12. The status of azimuthal velocity $g(\zeta)$ is enhanced with respect to fluid motion for the greater values of Maxwell nanofluid parameter β_1 which is clear from Fig. 13.

The axial velocity $h(\zeta)$ has increasing behaviors for the stretching parameter Ω_1 , porosity parameter λ_3 , Darcy Forchheimer parameter λ_4 , and Hall parameter m represented in Figs. 14, 15, 16, and 17 respectively. Figures 18 and 19 reflect on the decreasing effects of magnetic field parameter M and Maxwell nanofluid parameter β_1 for the axial velocity $h(\zeta)$ due to resistive behaviors.

Figure 20 is constructed to show the effect of stretching parameter Ω_1 on heat transfer $\theta(\zeta)$. On expansion of the disk, the temperature of the disk is reduced. The impact of Prandtl number Pr on temperature $\theta(\zeta)$ is given in Fig. 21. The figure shows a reducing behavior of heat transfer of Maxwell nanofluid by assuming different values of Prandtl number. In fact, for the high values of Prandtl number, the thermal diffusion is low; consequently, the fall in the heat transfer is occurred. In Figs. 22 and 23, it is observed that for the larger values of porosity and Darcy Forchheimer parameters (λ_3, λ_4), the expansion of temperature distribution $\theta(\zeta)$ is noted. The influence of thermal relaxation time parameter γ_3 on temperature $\theta(\zeta)$ is portrayed in Fig. 24. Both the thermal boundary layer and heat transfer are increased when the value of thermal relaxation time parameter γ_3 is increased. Physically, for the increasing values of γ_3 , the particles of fluid transfer more heat to adjacent particles. Figure 25 is plotted to show the effect of chemical reaction parameter k_c

on heat transfer $\theta(\zeta)$. For the higher values of k_c , the concentration is increased which increases the temperature. Figure 26 displays the behaviors of Maxwell nanofluid parameter β_1 and temperature profile $\theta(\zeta)$. The temperature boundary layer is decreased because more heat is required for the particles of Maxwell nanofluid due to its viscosity. The role of heat generation parameter β_2 is discussed in Fig. 27. Clearly, $\theta(\zeta)$ boosts up for the larger values of β_2 since the heat source/sink coefficient is increased. It is due to the fact that heat source/sink parameter β_2 creates an extra heat in the fluid. Figure 28 demonstrates the significance of Lewis number Le and heat transfer $\theta(\zeta)$. It is clear that temperature is an increasing function of Le . Figure 29 depicts the result that with higher estimates of Brownian motion parameter Nb , the non-dimensional temperature $\theta(\zeta)$ is increased. In fact, with the rise of Nb , Brownian motion of the fluid particles increases which generates resistance to the mass transport and increases the heat. Figure 30 indicates that temperature $\theta(\zeta)$ is improved with higher values of thermophoresis parameter Nt . It is due to the fact that greater values of Nt enhance the thermophoretic force which results in increment in heat transport.

Figure 31 captures the decrease of nanoparticles concentration $\phi(\zeta)$ due to stretching parameter Ω_1 . Figures 32 and 33 depict the features of Lewis number Le and chemical reaction parameter k_c respectively. It is recognized that for the higher estimations of Le and k_c , the nanoparticles concentration $\phi(\zeta)$ is enhanced. Figure 34 scrutinizes the influence of thermal relaxation time parameter γ_3 on nanoparticles concentration $\phi(\zeta)$. Physically, with rise of γ_3 , the increase in the relaxation time is occurred which decreases the mass transport. The characteristics of the non-dimensional solutal relaxation time parameter γ_4 and nanoparticles concentration profile $\phi(\zeta)$ are seen in Fig. 35. An increment in γ_4 causes high concentration between the fluid and nanoparticles. The reason is that with increasing values of solutal relaxation time parameter, the propagation

of concentration layers is enhanced which increases the mass transport. Similarly, in Fig. 36, it is noted that the nanoparticles concentration $\phi(\zeta)$ is maximum with increasing values of heat source/sink parameter β_2 because the extra heat provides more conversion to high concentration. Figures 37, 38, 39, 40, and 41 illustrate the effects of porosity parameter λ_3 , Darcy Forchheimer parameter λ_4 , Prandtl number Pr , Brownian motion parameter Nb , and thermophoresis parameter Nt respectively on nanoparticles concentration $\phi(\zeta)$. It is interpreted that $\phi(\zeta)$ attains maximum value for all these parameter estimations.

Figure 42 characterizes that the high values of Peclet number Pe lead to the reduction of gyrotactic microorganisms concentration $\chi(\zeta)$. It is clearly shown that the Peclet number Pe is very closely related to gyrotactic microorganisms concentration $\chi(\zeta)$ and nanoparticles concentration $\phi(\zeta)$ written in Eq. (15), i.e., $\chi'' - Lb\chi' - Pe(\chi'\phi' + \phi''(\gamma_5 + \chi)) = 0$. Similarly in Fig. 43, it is observed that the various values of bioconvection Lewis number Lb lead to the minimum value of gyrotactic microorganisms concentration $\chi(\zeta)$ which is also evident from Eq. (15). Figures 44, 45, and 46 are delineated to show the effects of Lewis number Le , chemical reaction parameter k_c , and the heat source/sink parameter β_2 on the gyrotactic microorganisms concentration $\chi(\zeta)$ respectively. Due to these parameters, the motion of gyrotactic microorganisms is decreased. Figure 47 indicates that the flow of gyrotactic microorganisms is enhanced by the increasing values of thermal relaxation time parameter γ_3 since time relaxation is provided. The result of microorganisms concentration difference parameter γ_5 is revealed through Fig. 48. Strong microorganisms concentration ultimately uplifts the $\chi(\zeta)$. In Fig. 49, $\chi(\zeta)$ is the increasing function of Brownian motion parameter Nb . Figure 50 indicates that the thermophoresis parameter Nt increases the microorganisms concentration $\chi(\zeta)$ profile.

Figure 51 elucidates the better performance of entropy generation rate $N_G(\zeta)$ due to temperature difference parameter γ_6 . The irreversibility of the present system is reduced. In Fig. 52, the result of increasing different values of diffusion constant parameter due to nanoparticles concentration γ_7 is indicated. Enhancement is observed due to γ_7 .

7 Conclusions

The Cattaneo-Christov heat and mass flux model is proposed taking into account the Maxwell nanofluid flow in a porous medium with gyrotactic microorganisms, entropy generation, and Hall current effect for the spinning disk. The homotopy analysis method (HAM) is applied for the computation of transformed equations resulting through similarity transformation. Heat and mass transfer are enhanced due to Cattaneo-Christov heat and mass flux model.

The main results are summarized as follows.

- (1) The radial velocity is reduced with the increasing values of Hall, Maxwell nanofluid, and magnetic field parameters while it is enhanced with porosity, Darcy Forchheimer, and stretching parameters.
- (2) The azimuthal velocity is reduced with the increasing values of stretching, porosity, and Darcy Forchheimer parameters while it is enhanced with the increasing values of Hall, Maxwell nanofluid, and magnetic field parameters.
- (3) The axial velocity is reduced with the increasing values of Maxwell nanofluid and magnetic field parameters while it is enhanced with porosity, Darcy Forchheimer, Hall, and stretching parameters.
- (4) The temperature is reduced with the increasing values of Prandtl number, stretching, Maxwell nanofluid, and chemical reaction parameters while it is enhanced with porosity, Darcy Forchheimer, thermal relaxation time, heat source/sink, Brownian motion, thermophoresis parameters, and Lewis number.
- (5) The nanoparticles concentration is reduced with the increasing values of stretching and thermal relaxation time parameters and Schmidt number while it is enhanced with Lewis and Prandtl numbers, chemical reaction, the porosity, Darcy Forchheimer, solutal relaxation time, heat source/sink, Brownian motion, and thermophoresis parameters.
- (6) The gyrotactic microorganisms concentration is reduced with the increasing values of Peclet, bioconvection Lewis and Lewis numbers, chemical reaction, and heat source/sink parameters while it is enhanced with the increasing values of thermal relaxation time, thermophoresis, gyrotactic microorganisms concentration difference, and Brownian motion parameters.
- (7) Entropy generation is reduced with temperature difference parameter while it is enhanced with values of diffusion constant parameter due to nanoparticles concentration.
- (8) Numerical comparison with published work provides a nice agreement.

Acknowledgements The authors are thankful to the respectable reviewers for their comments and suggestions. The authors acknowledge the financial support provided by the Center of Excellence in Theoretical and Computational Science (TaCS-CoE), KMUTT. Moreover, this research was supported by The Science, Research and Innovation Promotion Funding (TSRI) (Grant no. FRB650070/0168). This research block grants was managed under Rajamangala University of Technology Thanyaburi (FRB65E0633M.2). The authors extend their appreciation to the Deanship of Scientific Research at King Khalid University, Abha, Saudi Arabia for funding this work through research groups program under grant number RGP.1/248/43. The first author is thankful to the Higher Education Commission (HEC) Pakistan for funding through the SRGP project No. 10534.

Author contribution All authors contributed to the study conception and design. Conceptualization: Noor Saeed Khan, methodology: Poom Kumam, formal analysis: Usa Wannasingha Humphries, investigation: Wiyada Kumam, resources: Noor Saeed Khan, supervision: Taseer Muhammad, writing—original draft: Taseer Muhammad, writing—review and editing: Noor Saeed Khan.

Data availability Upon reasonable request to the corresponding author, the data will be provided.

Declarations

Conflict of interest The authors declare no competing interests.

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